

Don Freed

Addendum

(1)

1. RR fields

It is an abelian gauge field.

For string theory, maps

$$\Sigma^2 \longrightarrow \text{Minkowski}^{10}$$

quantize $\dots \rightarrow$ $\mathbb{Z}/2$ -graded Hilbert space

$$\mathcal{H} = \mathcal{H}^0 \oplus \mathcal{H}^1$$

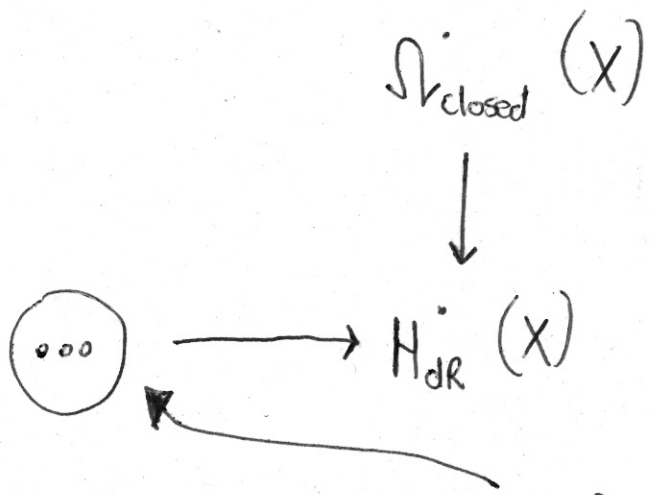
repⁿ of Poincaré.

You only see the massless fields at long distance...
diff. forms of high degree
... RR fields.

#

So they lie on Minkowski¹⁰.

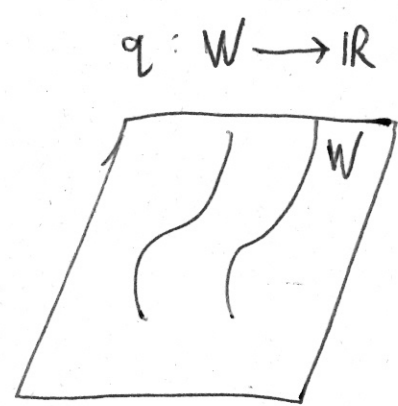
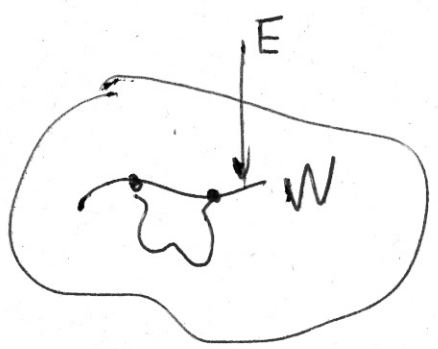
Put them on arbitrary X



For RR fields, want this to be $K^0(X)$

Why? • For anomaly cancellation reasons.

• Have to make 10d theory and 2d theory work.



$$\int_W \sqrt{\frac{\hat{A}(W)}{\hat{A}(V)}} \text{ch} E_{i^*} \rightarrow C$$

↑
RR field

plays the role of q

Maxwell

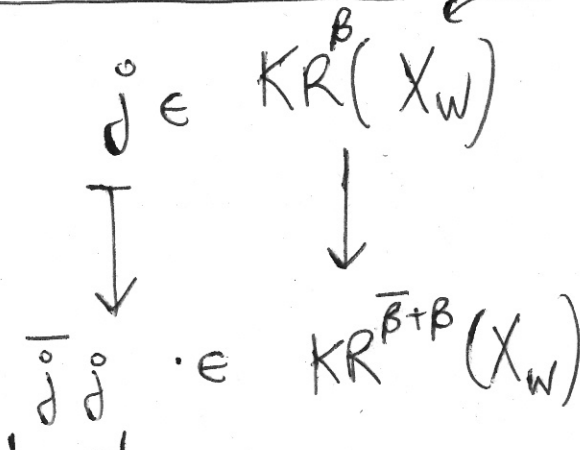
$$\int_W q i^* A$$

NS NS superstring background

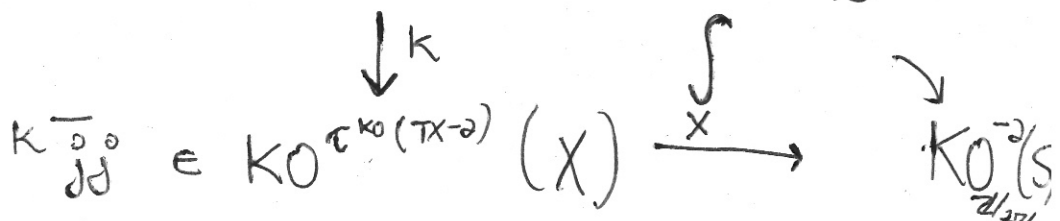
- ① X^{10} orbifold - metric, dilaton
- ② $\begin{matrix} (X_w)_\sigma \\ \downarrow \\ X \end{matrix}$ orientifold double cover
- ③ \checkmark
 β differential twisting of $KR(X_w)$
- ④ $k: R(\beta) \xrightarrow{\cong} \tau^{KO}(TX-2)$ iso of twistings of $KO(X)$.

What is this quadratic form?

the β -field is something which twists KR .



real structure, so lifts to $KO^{R(\beta)}(X)$.



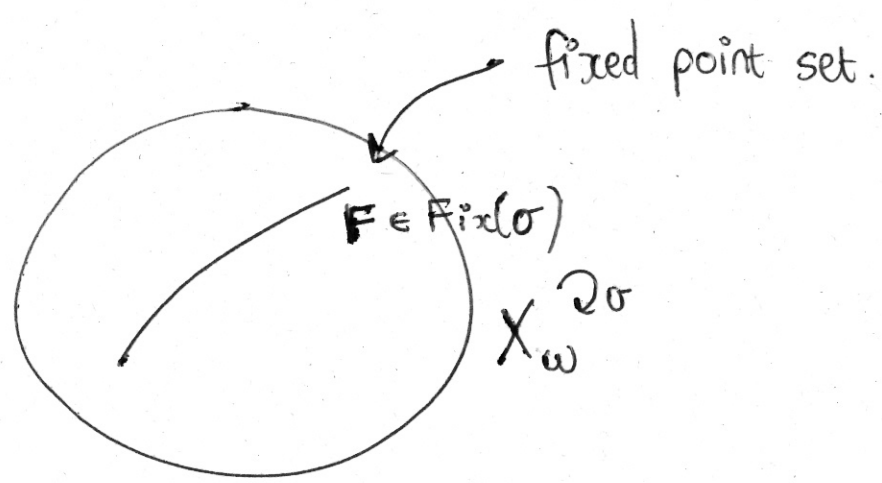
$$\int_{X/S} k \cdot \underbrace{j \cdot j}_{\text{real structure}}$$

take coefficient
→
of sign representation

$$H^2(S)$$

This is the input to the theory.

Now we can compute a formula for μ .



Use Atiyah-Singer localization theory in equivariant KO-theory.

We localize in this ring.

Wu class

$$\mu = i_* \left(\frac{\chi(F)}{\psi^{-1}(\text{Euler}^{KO}(U))} \right) \in KR^{\beta}[\frac{1}{2}](X)$$

Adams square localized

This is the story of the Ramond-Ramond field.

(3)

The cohomology theory needs to be self-dual

K theory ✓

bordism X

URS Maxwell theory from 1850

is secretly about gerbes with connection.

