

# Brownian loops and conformally invariant systems

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GTP Seminar

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جامعة نيويورك أبوظبي



# Outline

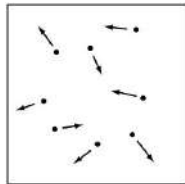
- 1 Conformal field theories
  - Statistical mechanics
  - Conformal invariance
  - Examples
- 2 Brownian loop soup
  - Primary operator spectrum
  - Correlation functions
- 3 Percolation
- 4 Overview and summary



## Statistical mechanics

Configuration space  $\Omega$  with Gibbs measure for  $\omega \in \Omega$

$$P(\omega) = \frac{1}{Z} e^{-\beta E(\omega)}$$



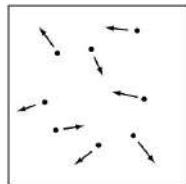
# Statistical mechanics

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Partition function

$$Z = \sum_{\omega \in \Omega} e^{-\beta E(\omega)}$$

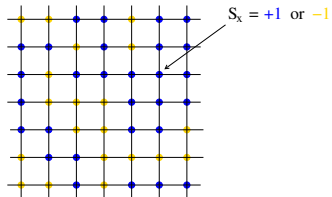


Correlation functions

$$\langle \mathcal{O}_1(x_1) \mathcal{O}_2(x_2) \dots \rangle = \frac{1}{Z} \sum_{\omega \in \Omega} \mathcal{O}_1(x_1) \mathcal{O}_2(x_2) \dots e^{-\beta E(\omega)}$$

## Example: Ising Model

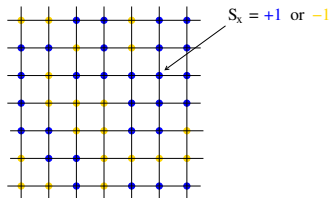
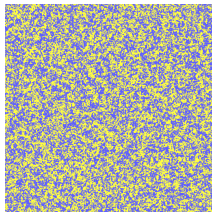
$$\Omega = \{-1, 1\}^D$$
$$E(\omega) = - \sum_{\langle i,j \rangle} S_i S_j$$



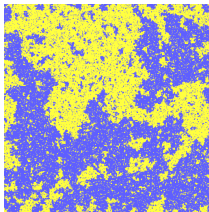
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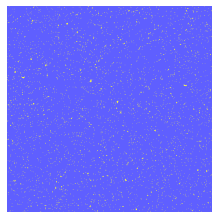
$$E(\omega) = - \sum_{\langle i,j \rangle} S_i S_j$$


 $\beta < \beta_{\text{critical}}$ 


Disordered phase

 $\beta_{\text{critical}}$ 


Critical point

 $\beta > \beta_{\text{critical}}$ 


Collective behavior

# Conformal transformations

Conformal field theories are quantum field theories that are invariant under angle-preserving transformations

$$\begin{aligned}x &\rightarrow x' = x + \varepsilon(x) \\g_{\mu\nu}(x) &\rightarrow g'_{\mu\nu}(x') = \Omega(x)g_{\mu\nu}(x)\end{aligned}$$

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Conformal transformations form the conformal group.

For  $d \geq 3$

$$SO(d+1, 1)$$

with  $\frac{1}{2}(d+1)(d+2)$  generators:  
translations, dilations, rotations, SCT



## Conformal transformations

$d = 2$  (Euclidean field theory) is special. Conformal transformations are given by Cauchy–Riemann equations.

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$$\bar{z} \rightarrow \bar{f}(\bar{z})$$

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generate conformal transformations.

The (quantum) generators form the Virasoro algebra

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12} n(n^2 - 1)\delta_{n+m,0}$$

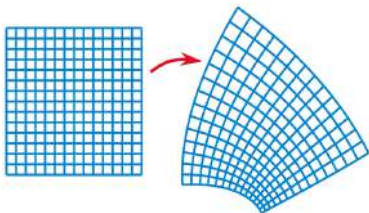
$$[L_n, \bar{L}_m] = 0$$

The global conformal transformations  $L_{-1}, L_0, L_1$  form a sub-algebra

# Conformal transformations

Any analytic function with nonzero derivative is a conformal map

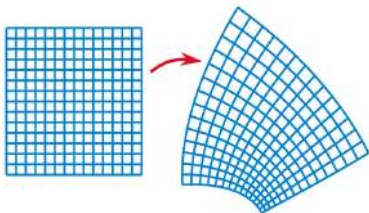
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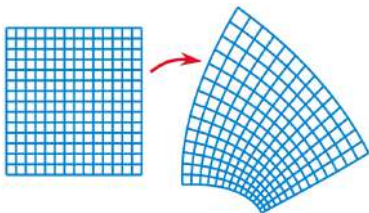
(Quasi-)primary operators behave like

$$\Phi(f(z), \bar{f}(\bar{z})) = \left( \frac{\partial f}{\partial z} \right)^{-\Delta} \left( \frac{\partial \bar{f}}{\partial \bar{z}} \right)^{-\bar{\Delta}} \Phi(z, \bar{z})$$

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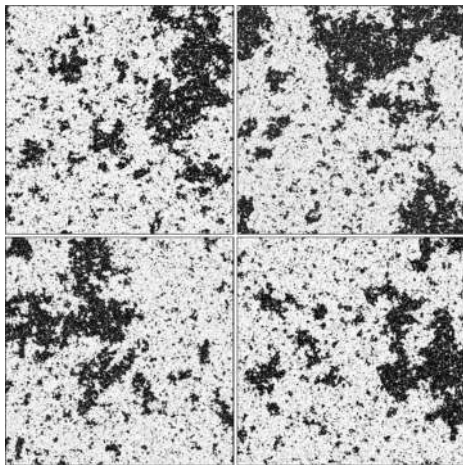


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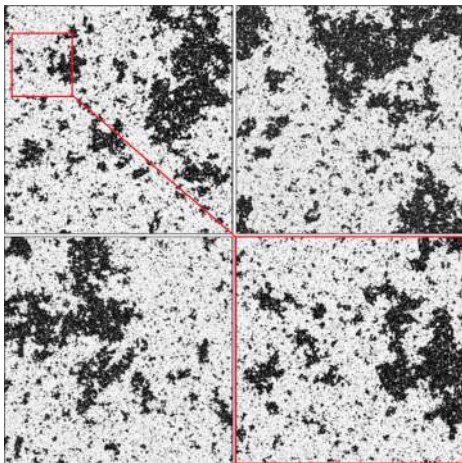
From here on  $\langle \Phi_1(z_1, \bar{z}_1) \dots \rangle \equiv \langle \Phi_1(z_1) \dots \rangle$

## Self-similarity: Ising model



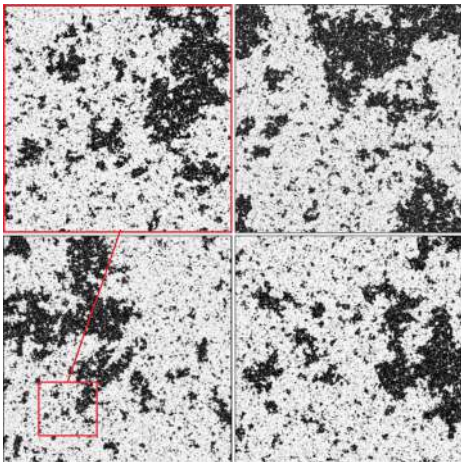
$$E = - \sum_{\langle i,j \rangle} S_i S_j, \quad S_i = \pm 1 \text{ at criticality}$$

## Self-similarity: Ising model

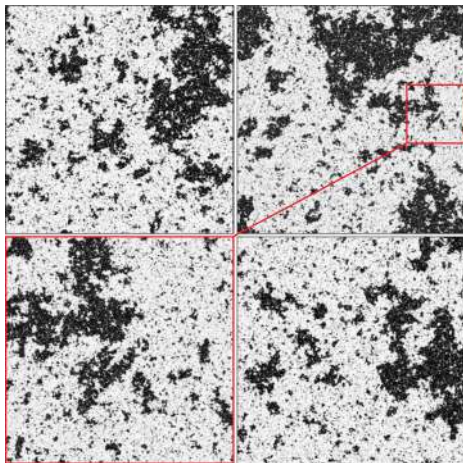




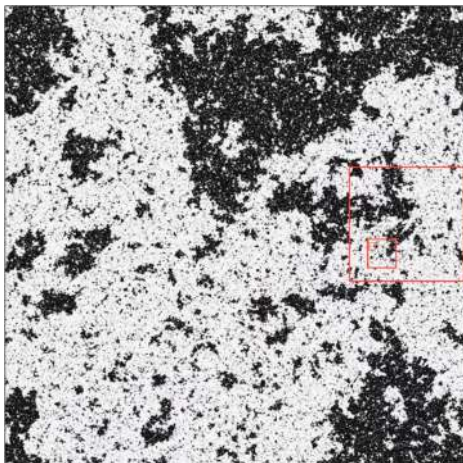
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In the continuum:  $\Phi(\lambda z) = \lambda^{-\Delta} \Phi(z)$

## Correlation functions

Every global conformal symmetry corresponds to one Ward identity:

$$\langle \Phi(z) \rangle$$

$$\langle \Phi(z_1) \Phi(z_2) \rangle$$

$$\langle \Phi_1(z_1) \Phi_2(z_2) \Phi_3(z_3) \rangle$$

are determined.

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The **four-point function**

$$\langle \Phi_1(z_1) \Phi_2(z_2) \Phi_3(z_3) \Phi_4(z_4) \rangle = G(x) \prod_{i < j}^4 z_{ij}^{c_{ij}} \bar{z}_{ij}^{\bar{c}_{ij}}$$

with  $z_{ij} = z_i - z_j$  and cross-ratio  $x = \frac{z_{12} z_{34}}{z_{13} z_{24}}$  contains non-trivial information

# CFTs appear in

- Statistical physics
- Condensed matter physics
- String theory
- AdS/CFT
- Stochastic systems



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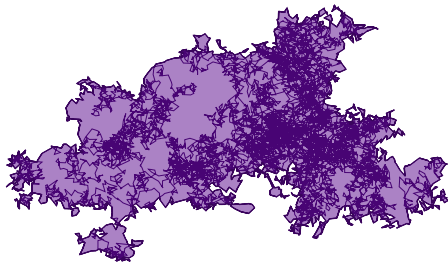
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Let's cook a Brownian loop soup!

# Brownian loop soup

The BLS is the Poissonian ensemble of Brownian loops in the plane

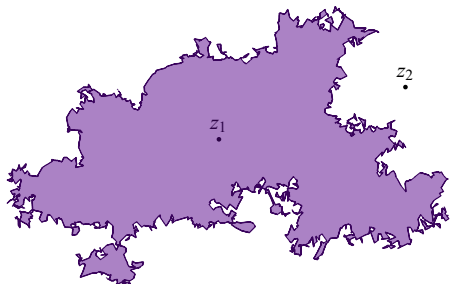
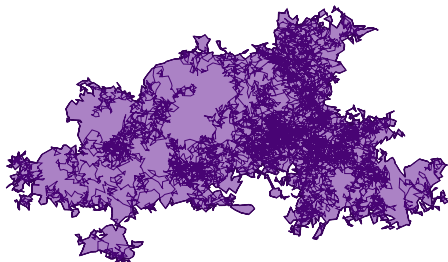


The outer boundary of a random walk is  
a self-avoiding random walk



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The layering number

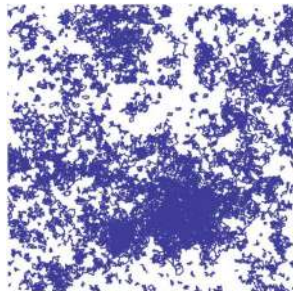
$$N(z_1) = 1$$

$$N(z_2) = 0$$

# Brownian loop soup

The BLS in domain  $D$  with intensity  $\lambda$

- is a Poisson point process
- of unrooted loops
- with unique measure  $\lambda\mu^{\text{loop}}$



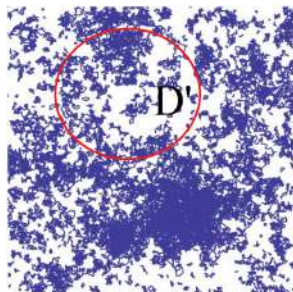
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The **restriction** of  $D$  to  $D' \in D$  is a BLS in  $D'$



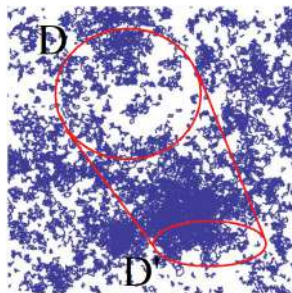
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The image of  $D \rightarrow D'$  under **conformal map** is BLS in  $D'$



## The layering vertex operators

The partition function

$$Z = \sum_{n=0}^{\infty} \frac{1}{n!} \left( \lambda \mu^{\text{loop}} \right)^n$$

Loop measure

$$\mu^{\text{loop}} = \int_D \int_0^{\infty} \frac{1}{2\pi t^2} \mu_{z,t}^{\text{br}} dt dA(z)$$

Central charge

$$c = 2\lambda$$

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Primary operators

$$e^{i\beta N(z)}$$

with  $N(z) = N_+(z) - N_-(z)$  have conformal dimension

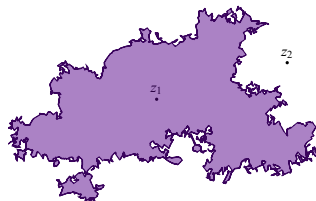
$$\Delta = \bar{\Delta} = \frac{\lambda}{10} (1 - \cos \beta)$$

## Correlation functions

Example: The two-point function

$$\left\langle e^{i\beta N(z_1)} e^{-i\beta N(z_2)} \right\rangle = \begin{cases} |z_1 - z_2|^{-4\Delta} & \text{CFTs} \\ \exp[-\lambda(1 - \cos\beta)(\alpha(z_1|z_2) + \alpha(z_2|z_1))] & \text{BLS} \end{cases}$$

$$\alpha(z_1|z_2) = \mu^{\text{loop}}(\text{loops covering } z_1 \text{ but not } z_2)$$

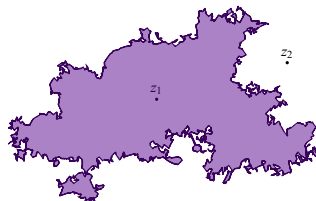


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Compute the four-point function

$$\left\langle \prod_{i=1}^4 e^{i\beta_i N(z_i)} \right\rangle = \exp[-\lambda(\text{sum of weights})]$$

These  $\alpha$ s appear in the  $O(n)$  model.

$n \rightarrow 0$  gives the ensemble of single, self-avoiding loops



# The $O(n)$ model

Random  $n$ -dimensional vectors:

- $n = 1$  Ising
- $n = 2$  XY (Berezinskii–Kosterlitz–Thouless)
- $n \rightarrow 0$  single, self-avoiding loops

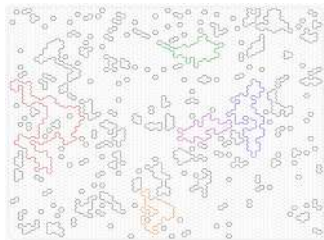
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$$\begin{aligned} Z(n) &= \int_{\mathbf{s}} d\mathbf{S} \prod_{\langle i,j \rangle} (1 + k\mathbf{S}_i \cdot \mathbf{S}_j) \\ &= \sum k^E n^L \\ &= 1 + n(\dots) + O(n^2) \end{aligned}$$

Sum over self-avoiding loop configurations  
of  $L$  loops and  $E$  edges

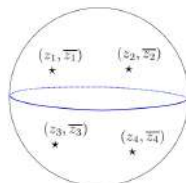
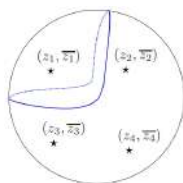
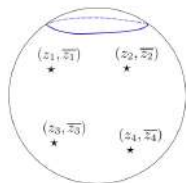


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Random  $n$ -dimensional vectors:

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- $n \rightarrow 0$  single, self-avoiding loops



The BLS for  $\lambda \rightarrow 0$  becomes the  $O(0)$  model. Study the critical  $O(n)$  model as a CFT!

## A four-point function

All weights are determined

$$\alpha_{\mathbb{S}}(z_1 | z_2, z_3, z_4) = \frac{1}{5} \left( \log \left| \frac{z_{12} z_{14}}{z_{24}} \right| + A(x) \right)$$

$$\alpha_{\mathbb{S}}(z_1, z_2 | z_3, z_4) = -\frac{1}{5} (\log |x| + A(x))$$

$$\vdots$$

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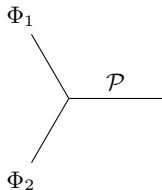
$$\left\langle \prod_{i=1}^4 e^{i\beta_i N(z_i)} \right\rangle = \exp \left[ \tilde{\Delta} A(x) \right] \prod_{i < j}^4 |z_{ij}|^{c_{ij}}$$

$$A(x) = \frac{1}{4} [x {}_3F_2(x) + \bar{x} {}_3F_2(\bar{x})] - \text{const} \cdot |x(1-x)|^{\frac{2}{3}} |{}_2F_1(x)|^2$$

## Conformal block expansion

Operator Product Expansion (fusion rules)

$$\Phi_1(z + \varepsilon)\Phi_2(z) = \sum_{\mathcal{P}} \varepsilon^{\delta_{\mathcal{P}}} \bar{\varepsilon}^{\bar{\delta}_{\mathcal{P}}} C_{12}^{\mathcal{P}} \mathcal{P}(z)$$

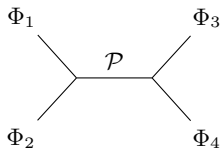
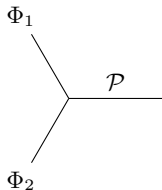


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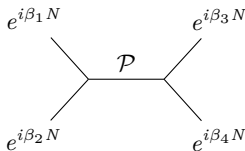
Apply to four-point function



# Conformal block expansion

The Virasoro conformal block expansion

$$G(x) = \sum_{\mathcal{P}} C_{34}^{\mathcal{P}} C_{12}^{\mathcal{P}} \mathcal{F}(\mathcal{P}|x) \bar{\mathcal{F}}(\mathcal{P}|\bar{x})$$



$$\Delta_{\mathcal{P}} = \Delta_{12} + \frac{p}{3}$$

$$\bar{\Delta}_{\mathcal{P}} = \Delta_{12} + \frac{p'}{3}$$

$$\Delta - \bar{\Delta} \in \mathbb{Z}$$

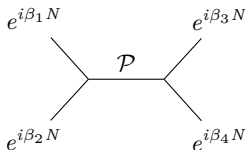
$$\Delta_{12} = \frac{\lambda}{10} (1 - \cos(\beta_1 + \beta_2))$$



# Conformal block expansion

The Virasoro conformal block expansion

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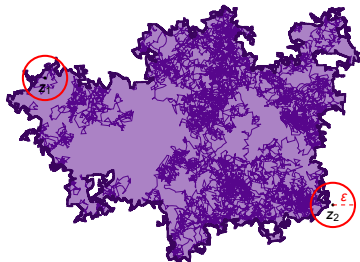
contains information about the entire CFT

- Central charge  $c$
- Primary operator spectrum  $\Delta_{\mathcal{P}}$
- Three-point function coefficients  $C_{ij}^{\mathcal{P}}$

# Edge counting operators

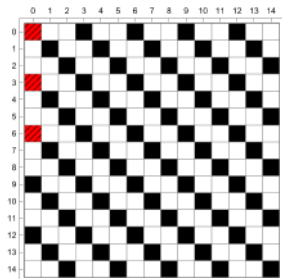
What are the new operators in terms of loops? Define

$E_\epsilon(z) :=$  Centered number of loops whose outer boundaries come  $\epsilon$ -close to  $z$



$$\Delta_\epsilon = \bar{\Delta}_\epsilon = \frac{1}{3}$$

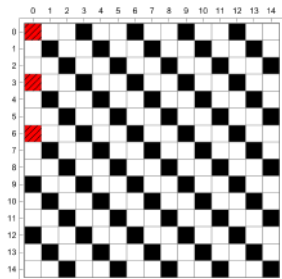
## Extended Symmetry?



$$(\Delta_{\mathcal{P}}, \bar{\Delta}_{\mathcal{P}}) = (p/3, p'/3)$$

$$\Delta_{\mathcal{P}} - \bar{\Delta}_{\mathcal{P}} = \mathbb{Z}$$

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$$T(z) = (L_{-2}\mathbb{1})(z)$$

$$J(z) = (J_{-1}\mathbb{1})(z)$$

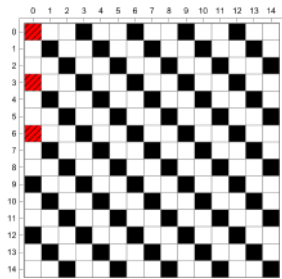
$$[J] = (1, 0)$$

$$\psi = J_{-m_1} J_{-m_2} \dots \phi$$

$$L_0 \phi = \Delta \phi$$

$$L_0 \psi = (\Delta + m_1 + m_2 + \dots) \psi$$

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$$\Delta_{\mathcal{P}} - \bar{\Delta}_{\mathcal{P}} = \mathbb{Z}$$

$$T(z)\phi_i(w) = \frac{\Delta_i}{(z-w)^2}\phi_i(w) + \frac{1}{z-w}\partial_w\phi_i(w)$$

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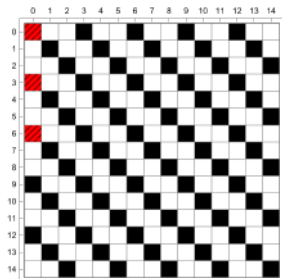
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$$\Delta_{\mathcal{P}} - \bar{\Delta}_{\mathcal{P}} = \mathbb{Z}$$

$$T(z)\phi_i(w) = \frac{\Delta_i}{(z-w)^2}\phi_i(w) + \frac{1}{z-w}\partial_w\phi_i(w)$$

$$J(z)\phi_i(w) = \frac{t_i}{z-w}\phi_i(w)?$$

$$T(z) = (L_{-2}\mathbb{1})(z)$$

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$$\psi = J_{-m_1}J_{-m_2}\dots\phi$$

$$L_0\phi = \Delta\phi$$

$$L_0\psi = (\Delta + m_1 + m_2 + \dots)\psi$$

## Some comments

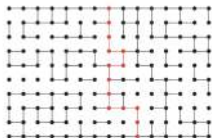
- The BLS cannot be unitary, at least for  $c < 1$
- We expect unitarity for  $\lambda = 1/2$  or  $c = 1$  (free Boson)
- The BLS can be generalized to domains with boundary
- What is the full symmetry algebra / operator spectrum? (Kac-Moody)
- Where does the BLS lie in the class of CFTs? (minimal models, logarithmic CFTs)



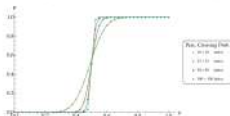
## Percolation

J. Phys. A: Math. Theor. 49 (2017) 084305

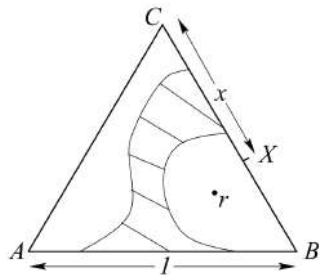
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**Figure 1.** A bond percolation sample on a square lattice inside a rectangle. This sample illustrates a crossing event: the collection of all activated bonds (black) contains a connected path (red) that joins the bottom side of the rectangle to the top side.



**Figure 2.** Crossing probability  $\bar{P}$  as a function of bond activation probability  $p$  for percolation on an  $M \times M$  square lattice in a square ( $R=1$ ). As  $M \rightarrow \infty$ ,  $\bar{P}$  approaches a step function (—) that jumps at the critical probability  $p_c^{sq} = 0.5$ .



Cardy's connection formula

$$P(\text{a cluster connects } AB \text{ and } CX) = x$$



## Conformal probability

Modern probability theory studies **random fractal** objects that are **conformally invariant** in distribution. (Continuum scaling limit at **criticality** of percolation, Ising, Potts,  $O(n)$ , loop-erased random walk, self-avoiding walk, ...)

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Schramm-Loewner evolution  
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Conformal loop ensembles  
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Brownian loop soup  
(Lawler, Werner)



Gaussian free field and GMC  
(Mandelbrot; Kahane, Peyrière)

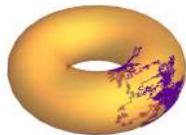


# Summary

Conformal probability provides a link between 2D statistical mechanics and conformal field theory (rigorous approach to CFT).

The Brownian loop soup

- allows a microscopic interpretation of the field theory objects
- has connections to the free Boson,  $O(n)$  model, SLE, CLE, ...
- is a rich model with new features



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Based on

arXiv:1912.00973, 2109.12116, 2112.00074, TBA  
with Federico Camia, Alberto Gandolfi, Matthew Kleban



# Thanks!