

Holographic QCD

Status and perspectives for the future

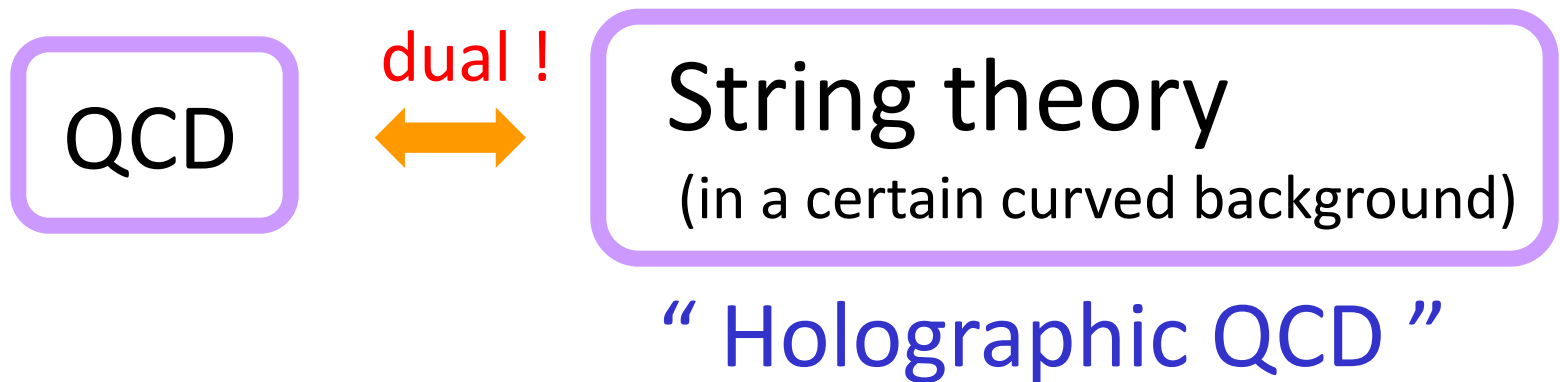
Shigeki Sugimoto (Kavli IPMU)

Claim :

**String theory can be
a theory of hadrons**

The string scale here is around 1 GeV.

I'm NOT saying QCD is wrong.



“Gauge/String duality”

[Maldacena 1997, ...]

How can we realize QCD in string theory?

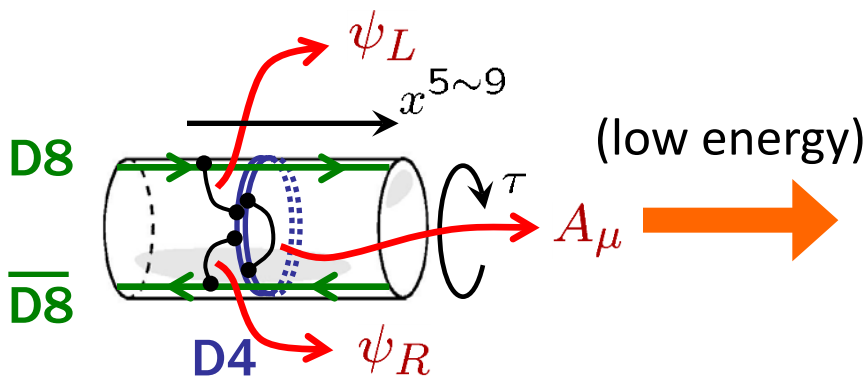
(1+4) dim
D-brane

(1+8) dim
D-brane

		x^0	x^1	x^2	x^3	τ	x^5	x^6	x^7	x^8	x^9
D4	$\times N_c$	○	○	○	○	○	-	-	-	-	-
D8- $\overline{\text{D8}}$	$\times N_f$	○	○	○	○	-	○	○	○	○	○

$\underbrace{\hspace{10em}}_{4 \text{ dim}}$
 $\curvearrowright S^1$

[Sakai-S.S. 2004]



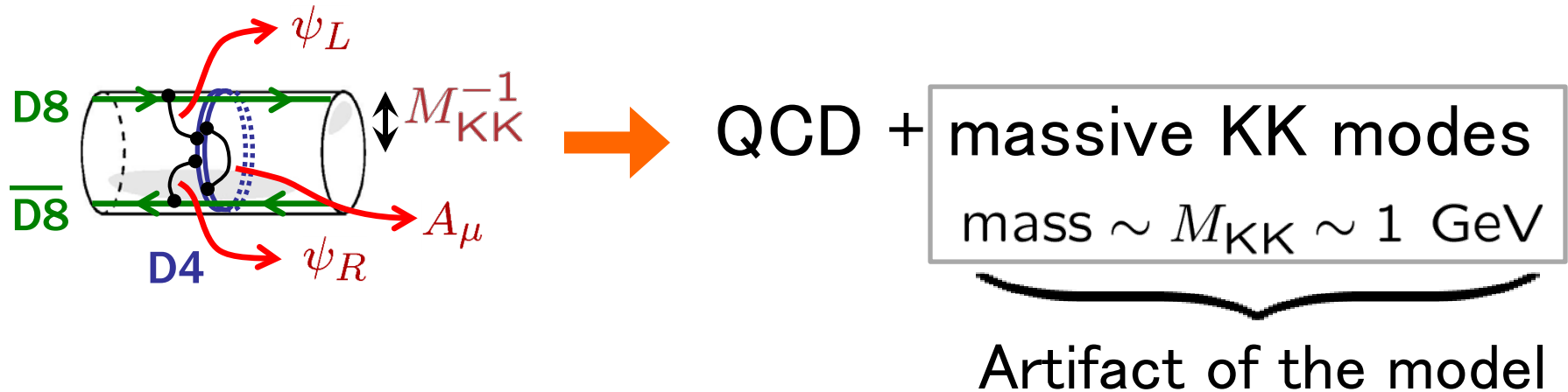
A_μ : gluon

$\begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$: quark $\times N_f$

$SU(N_c)$ QCD

with N_f massless quarks

Comments



- We try to get rid of the states that contain these unwanted modes as much as possible using a symmetry argument.
- These artifacts may contribute in loops. We simply hope the contributions are not large.
 - cf) quench approximation in lattice QCD

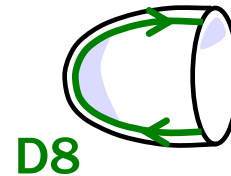
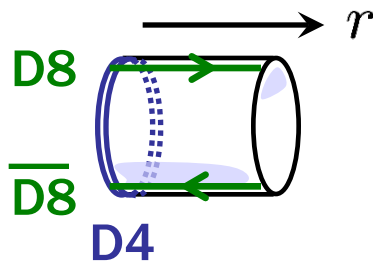
Holographic QCD

[Witten 1998]

replace D4
with curved
background

N_c D4-brane on S^1
+ N_f D8- $\overline{\text{D8}}$ pairs

String theory in
the D4 background
+ N_f D8-branes



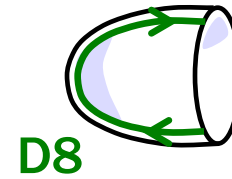
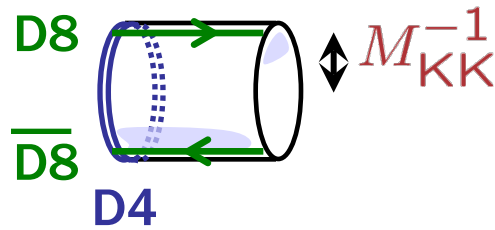
(low energy)

$SU(N_c)$ QCD
with N_f massless quarks

Open + closed string theory
in this background

↔
dual

“ Holographic QCD ”



(low energy)



$SU(N_c)$ QCD
with N_f massless quarks

\longleftrightarrow
dual

Open + closed string theory
in this background



parameters :

$$\left\{ \begin{array}{l} M_{KK} \sim \text{cut off scale} \\ \lambda = g_{YM}^2 N_c \\ \text{'t Hooft coupling} \end{array} \right.$$

$$\left\{ \begin{array}{l} \text{string coupling} \propto 1/N_c \\ \text{string length} \propto \lambda^{-1/2} \end{array} \right.$$



good description for

$$\left\{ \begin{array}{l} \text{large } N_c \\ \text{large } \lambda \longleftrightarrow \text{low energy} \end{array} \right.$$

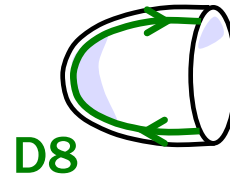
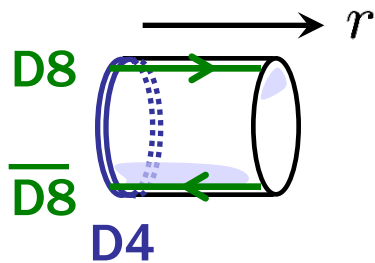
Chiral Sym Breaking

replace D4
with curved
background



N_c D4-brane on S^1
+ N_f D8- $\overline{\text{D8}}$ pairs

String theory in
the D4 background
+ N_f D8-branes



$$U(N_f)_L \times U(N_f)_R$$

\updownarrow
 D8

 \updownarrow
 $\overline{\text{D8}}$

$$U(N_f)$$

\updownarrow
 Connected D8

As expected in QCD!

Do we have hadrons?

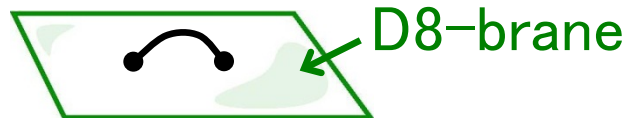
Particles in this system

- Closed strings



↔ glueballs

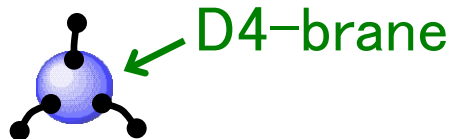
- Open strings on D8



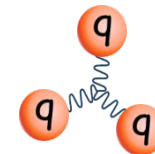
↔ mesons



- D4 wrapped on S^4



↔ baryons



N_c strings are attached

Surprise :

A lot of properties of hadrons can be extracted using this new description!

Caution :

- We have only estimated the leading terms in the $1/N_c$ and $1/\lambda$ expansions.
- M_{KK} is around 1 GeV.
- Quark masses are neglected.

We should not expect too much.
But, don't be too pessimistic!

- Meson effective theory is written as a 5 dim $U(N_f)$ YM-CS theory in a curved background.

$$\leftarrow (x^{0\sim 3}, z)$$

$$S_{5\text{dim}} \simeq S_{\text{YM}} + S_{\text{CS}}$$

$$S_{\text{YM}} = \kappa \int d^4x dz \text{Tr} \left(\frac{1}{2} h(z) F_{\mu\nu}^2 + k(z) F_{\mu z}^2 \right) \quad S_{\text{CS}} = \frac{N_c}{24\pi^2} \int_5 \omega_5(A)$$

$k(z) = 1 + z^2$ CS5-form
 $h(z) = (1 + z^2)^{-1/3}$

$\mu, \nu = 0 \sim 3$

$$A_\mu(x^\mu, z) = \sum_{n \geq 1} B_\mu^{(n)}(x^\mu) \psi_n(z)$$

$$A_z(x^\mu, z) = \sum_{n \geq 0} \varphi^{(n)}(x^\mu) \phi_n(z)$$

complete sets of functions of z

$$\varphi^{(0)} \sim \text{pion} \quad B_\mu^{(1)} \sim \rho \text{ meson} \quad B_\mu^{(2)} \sim a_1 \text{ meson} \quad \dots$$



$$S_{5\text{dim}}(A) = S_{4\text{dim}}(\pi, \rho, a_1, \rho', a'_1, \dots)$$

traditional meson effective action

● A lot of models are reproduced without making any phenomenological assumptions!

next slide

- Skyrme model [Skyrme 1961]
- Vector meson dominance [Sakurai 1960, Gell-Mann-Zachariasen 1961]
- Gell-Mann Sharp Wagner model [Gell-Mann –Sharpe-Wagner 1962]
- Hidden local symmetry [Bando-Kugo-Uehara-Yamawaki-Yanagida 1985]
- Son-Stephanov's bottom up holographic model [Son-Stephanov 2003]

● masses and couplings roughly agree with experimental data.

[Sakai-S.S. 2004, 2005]

mass	ρ	a_1	ρ'
exp.(MeV)	776	1230	1465
our model	[776]	1189	1607
ratio	[1]	1.03	0.911

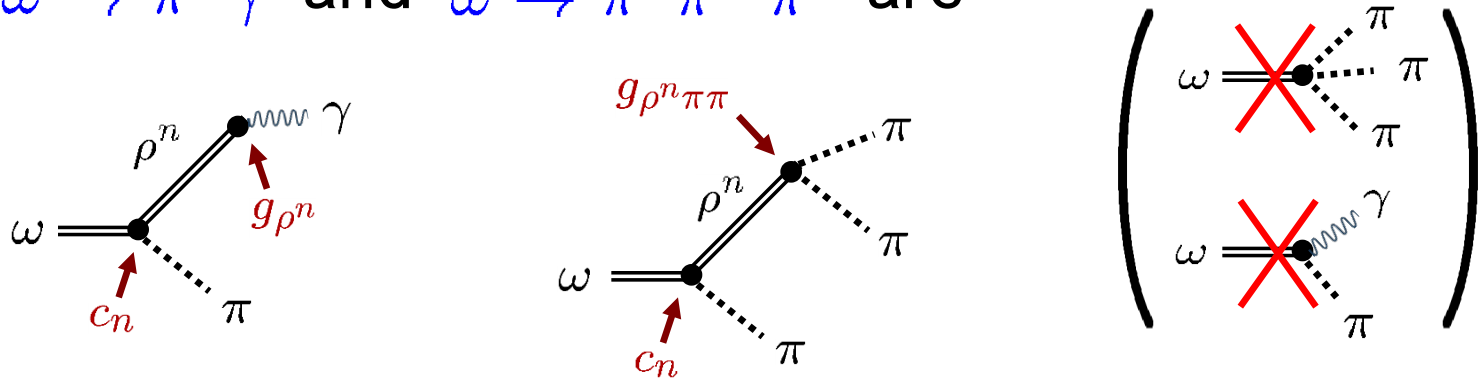
input

(to fix M_{KK}, λ)

coupling	our model	experiment
f_π	[92.4 MeV]	92.4 MeV
L_1	0.584×10^{-3}	$(0.1 \sim 0.7) \times 10^{-3}$
L_2	1.17×10^{-3}	$(1.1 \sim 1.7) \times 10^{-3}$
L_3	-3.51×10^{-3}	$-(2.4 \sim 4.6) \times 10^{-3}$
L_9	8.74×10^{-3}	$(6.2 \sim 7.6) \times 10^{-3}$
L_{10}	-8.74×10^{-3}	$-(4.8 \sim 6.3) \times 10^{-3}$
$g_{\rho\pi\pi}$	4.81	5.99
g_ρ	0.164 GeV ²	0.121 GeV ²
$g_{a_1\rho\pi}$	4.63 GeV	2.8 ~ 4.2 GeV

● ω meson decay ($\omega \rightarrow \pi^0 \gamma$ and $\omega \rightarrow \pi^0 \pi^+ \pi^-$)

- Our model predicts that the relevant diagrams for $\omega \rightarrow \pi^0 \gamma$ and $\omega \rightarrow \pi^0 \pi^+ \pi^-$ are



➔ Exactly the same as the **GSW model** !

[Gell-Mann -Sharp-Wagner 1962]

- Furthermore, we find

$$\Gamma(\omega \rightarrow \pi^0 \gamma) = \frac{N_c^2}{3} \frac{\alpha}{64\pi^4 f_\pi^2} \left(\sum_{n=1}^{\infty} \frac{c_n g_{\rho^n}}{m_{\rho^n}^2} \right)^2 |\mathbf{p}_\pi|^3 = \frac{N_c^2}{3} \frac{\alpha}{64\pi^4 f_\pi^2} g_{\rho^n \pi \pi}^2 |\mathbf{p}_\pi|^3$$

➔ reproduces the proposal given by Fujiwara et al !

[Fujiwara-Kugo-Terao-Uehara-Yamawaki 1985]

- Other mesons, including higher spin mesons, are obtained as excited open string states.

[Imoto-Sakai-S.S. 2010]



- 1st excited states**

→ $a_2(1320)$, $b_1(1235)$, $\pi(1300)$, $a_0(1450)$, ...

- 2nd excited states**

→ $\rho_3(1690)$, $\pi_2(1670)$, ...

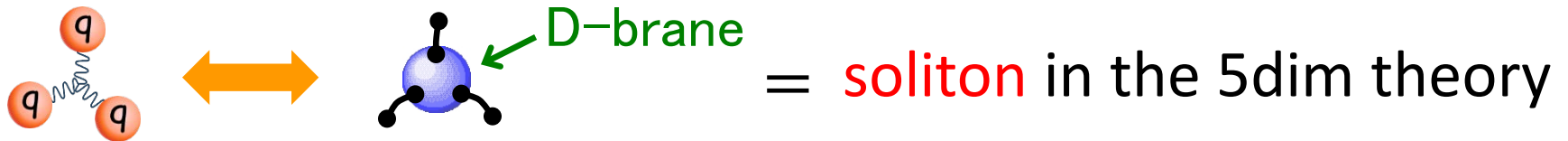
- The existence of isovector mesons with 2^{--} , 1^{+-} , 0^{++} around 1700 MeV is suggested.

- No good candidate states for $a_0(980)$, $\pi_1(1400)$

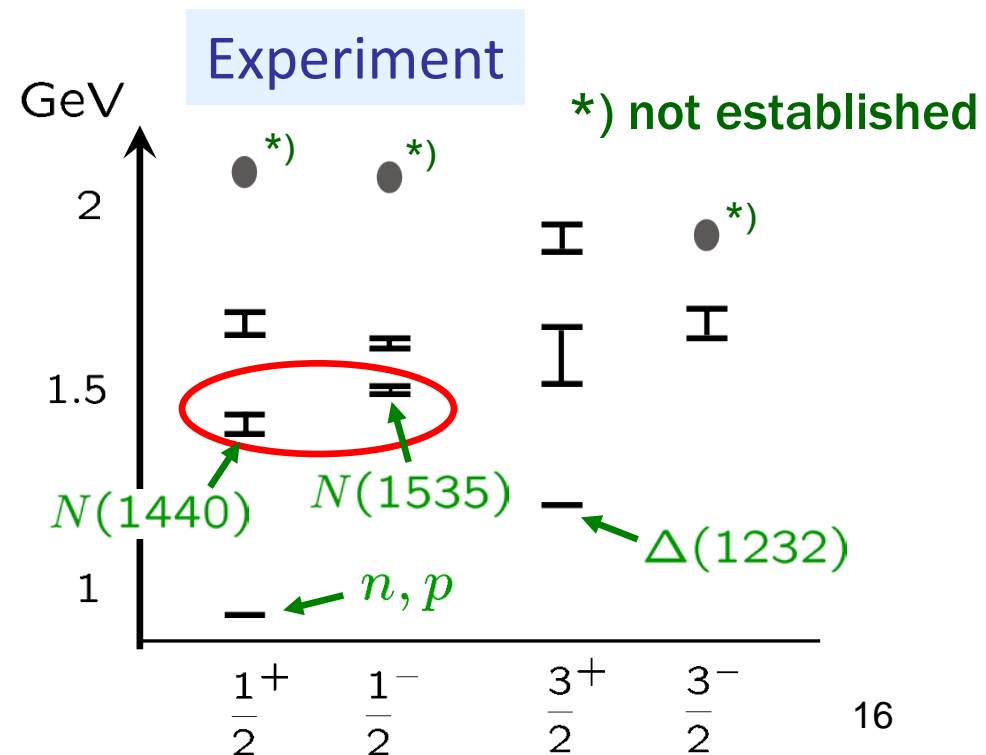
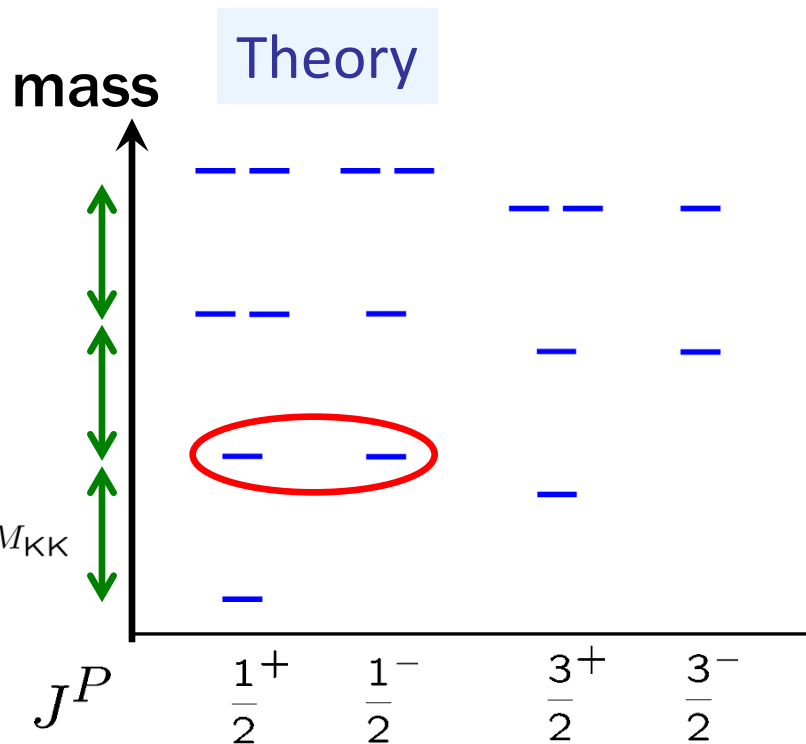
→ Suggesting that they are 4 quark states.

Baryon spectrum

[Hata-Yamato-Sakai-S.S. 2007]



We can analyze the spectrum, magnetic moments, charge radii by quantizing this soliton cf) Skyrme model



● Properties of nucleons

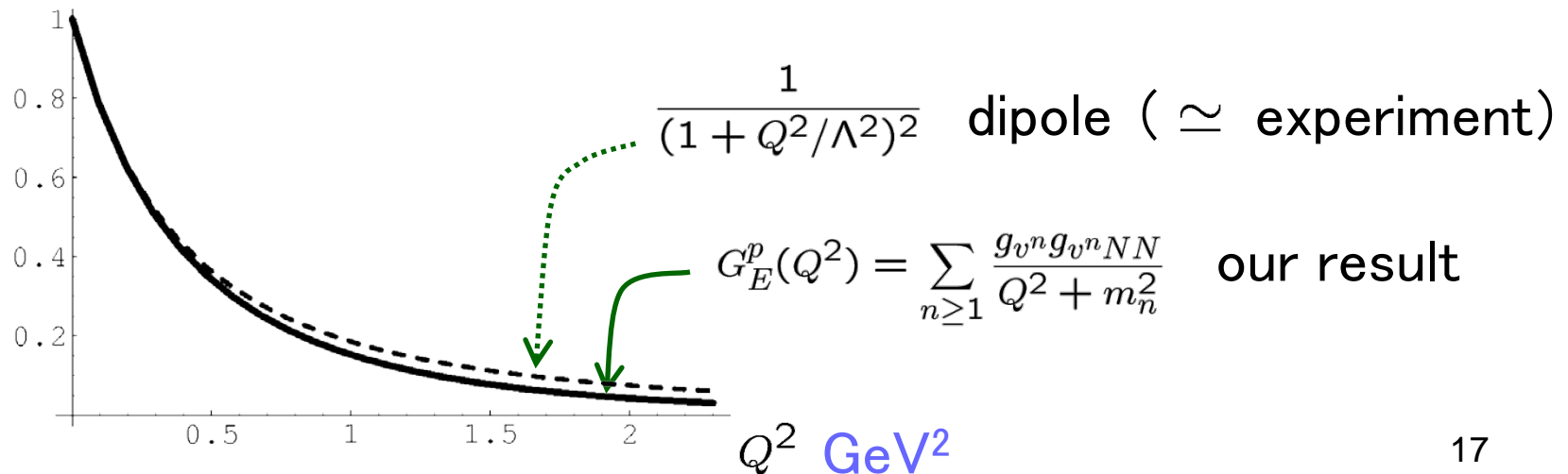
[Hashimoto-Sakai-S.S. 2008]

	our result	exp.
$\langle r^2 \rangle_{I=0}^{1/2}$	0.74 fm	0.81 fm
$\langle r^2 \rangle_{I=1}^{1/2}$	0.74 fm	0.94 fm
$\langle r^2 \rangle_A^{1/2}$	0.54 fm	0.67 fm
$g_{I=0}$	1.7	1.8
$g_{I=1}$	7.0	9.4
g_A	0.73	1.3

[See also,

Hong-Rho-Yee-Yi 2007,
Hata-Murata-Yamato 2008,
Kim-Zahed 2008]

● nucleon ele-mag form factor



Summary of Status

● Qualitative features

Confinement, Chiral symmetry breaking, ...
understood from the geometry of the background.

● Spectrum of hadrons

glueballs, mesons and baryons

↔ closed, open strings and D-branes

Comparison with experimental data is encouraging.

● Structure of interaction

Many old phenomenological models are reproduced.

A lot of couplings are calculated and compared
with experimental data.

Perspectives

- More accurate calculation

$1/\lambda$, $1/N_c$ corrections

$\longleftrightarrow \alpha'$, loop corrections in string theory

- Connection to perturbative QCD

Go beyond $1/\lambda$ expansion

Need some clever ideas

- Phase structure

Many works including T , μ , E , B , \dots

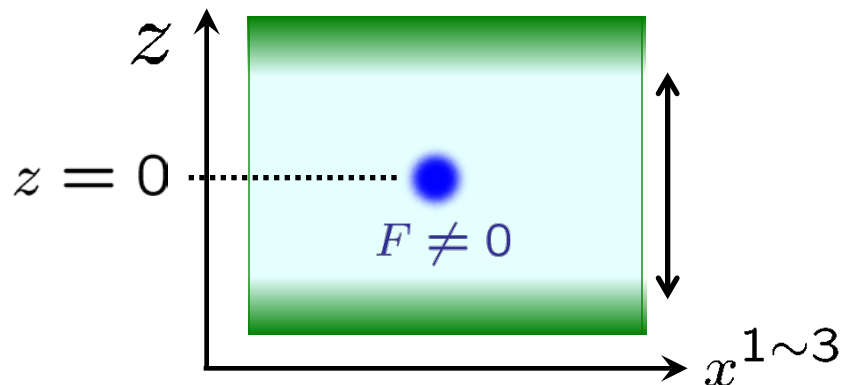
Actively developing area of research.

Large T and/or density are more challenging.

● What happens if we have many baryons?

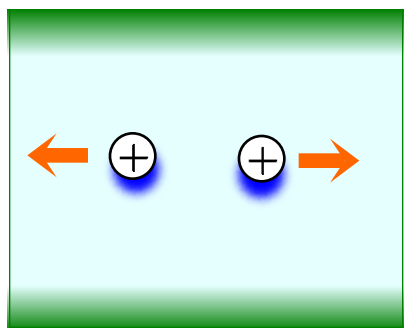
Baryon = **soliton** in 5dim $U(N_f)$ theory

↖ $(x^{0\sim 3}, z)$



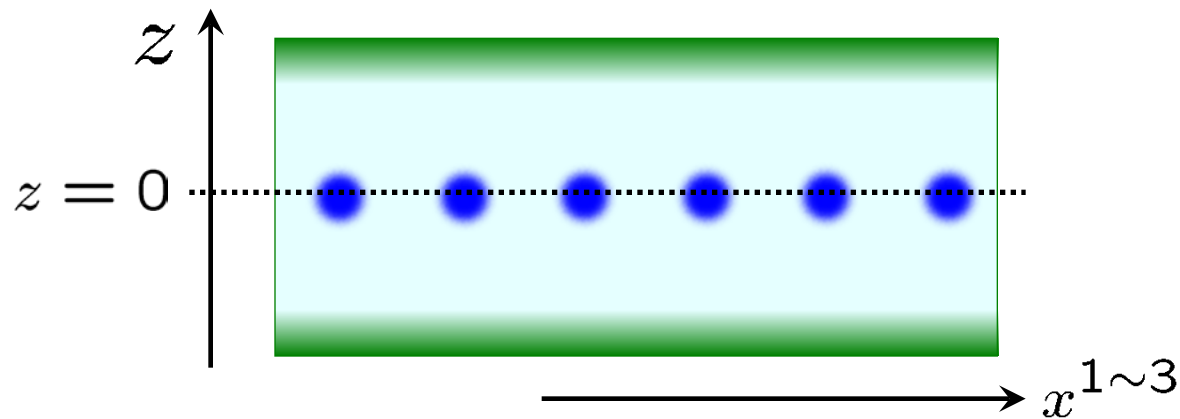
- Baryon # : $N_B = \frac{1}{8\pi^2} \int \text{Tr}(F \wedge F)$
- Size of the soliton is small
- Behaves like a point charge
- Soliton wants to be near $z=0$

- Repulsive at short distance mainly because of $U(1)$ charge.

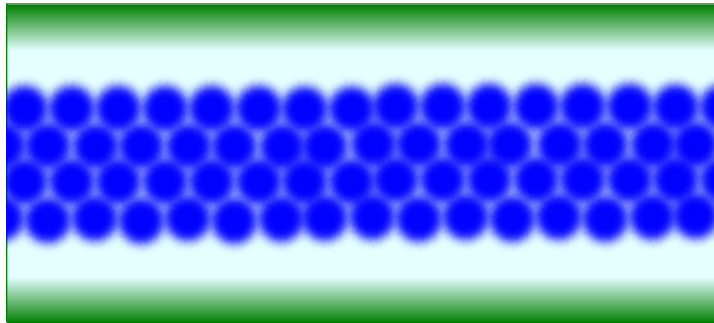


$U(1)$ part $\ni \omega$ -meson

- Finite density



- higher density



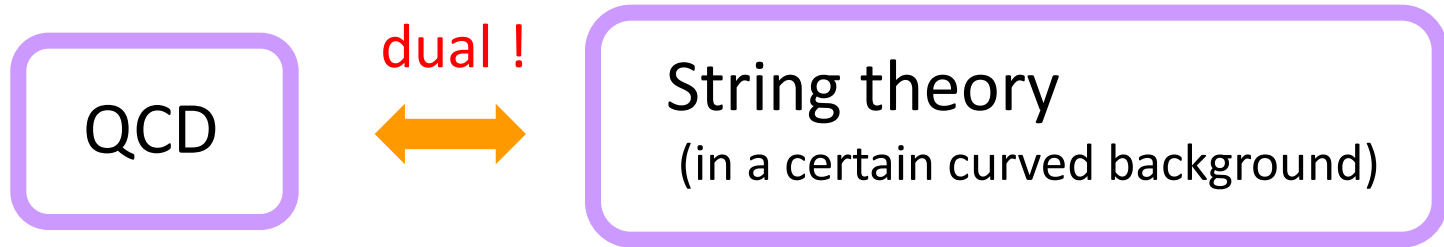
[Rozali-Shieh-Van Raamsdonk
-Wu 2007, Kaplunovsky-Melnikov
-Sonnenschein 2012, de Boer
-Chowdhury-Heller-Jankowski 2012]

Quarkyonic phase?

Interpretation is not clear to me yet.

Conclusion :

String theory can be
a theory of hadrons



“ Holographic QCD ”

