

# I Factorization algebras in perturbative quantum field theory

## By Kevin Costello (Northwestern) May 2009

### I.1 Deformation quantization (work in progress with Owen Guilliam).

Classical mechanics: described by a Poisson algebra  $(A, \{\cdot, \cdot\})$ . In deformation quantization, want to replace  $A$  by an associative product  $A[[\hbar]]$ , such that if  $f, g \in A$ ,  $\{f, g\} = \lim_{\hbar \rightarrow 0} \frac{1}{\hbar} [f, g]$ .

Want an analogy in QFT:

- 1) Describe classical and quantum algebraic structures.
- 2) Show that we can get classical structure from classical field theory.
- 3) State a "quantization" theorem.

Let  $M$  be a manifold. A *factorization algebra* on  $M$  can be described as follows. Let  $B(M)$  be the space of balls in  $M$  and  $B_n(M)$  the collection of  $n$  disjoint unions of balls  $B_1, \dots, B_n \in B(M)$  embedded into  $B_{n+1}$  (like in the definition of the little cubes operad).

A factorization algebra is a vector bundle  $V$  on  $B(M)$ , together with maps

$$B(M) \xleftarrow{p} B_n(M) \xrightarrow{q} B(M)^n, \quad q^*(V^{\otimes n}) \rightarrow p^*V$$

satisfying some evident compatibility:  $V(B_1) \otimes V(B_2) \rightarrow V(B_3)$ . The two maps

$$V^{\otimes \text{three balls}} \rightarrow V^{\otimes \text{two intermediate balls}} \rightarrow V^{\text{outer ball}}$$

and  $V^{\otimes \text{three balls}} \rightarrow V^{\text{outer ball}}$  commute.

$\rightsquigarrow$  Close relation of  $E_n$ -algebras for  $\dim(M) = n$ . Say a top. factorization algebra in  $M$  is this structure when  $V$  is a locally constant sheaf and the maps are morphisms of locally constant sheafs.

**Theorem I.1.** *An  $E_n$ -algebra yields a top. factorization algebra on any framed manifold  $M$  with  $n = \dim(M)$ .*

A factorization algebra in  $\mathbb{C}$ , where everything is holomorphic an invariant under  $\text{Aff}(\mathbb{C})$ :

$$V(\{z \in \mathbb{C} : |z| < 1\}) =: W \quad \text{a vector space}$$

and the factorization algebra gives a map  $m_z : W \otimes W \rightarrow W$ , depending holomorphically on  $z$  in an annulus. Thus  $m_z \sim \sum_{k \in \mathbb{Z}} \varphi_k z^k$ , where  $\varphi_k \in$  some completion of  $\text{Hom}(W, W) \rightsquigarrow$  reminiscent of the operator product expansion in vertex algebras.

**Claim:** Factorization algebras on  $M$  encode structure one expects from a quantum field theory on  $M$ . This is motivated by the 2-dim. holomorphic setting, which fits with the known picture. In one dimension, this reduces to an associative algebra, the algebra of observables of quantum mechanics. In any dimension, one can construct a factorization algebra using perturbative quantum field theory.

Factorization algebras are symmetric monoidal algebras. A classical fact. algebra is a commutative algebra in this category. Suppose we have a classical field theory. For instance, take  $M$  compact riemannian, fields to be  $C^\infty(M, \mathbb{R})$  and the action

$$S(\varphi) = \int_M \varphi \Delta \varphi + \varphi^3$$

Let  $EL$  be the sheaf of solutions to the Euler-Lagrange equations  $2\Delta\varphi + 3\varphi^2 = 0$ . If  $B \subseteq M$ , let  $\mathcal{O}(EL(B))$  be functions on  $EL(B)$ . This is a commutative fact. algebra  $EL(B_3) \rightarrow EL(B_1) \times EL(B_2)$  and applying  $\mathcal{O}$  to this gives a fact. algebra.

**Claim:** The  $E_\infty$ -algebra "wants" to become  $E_0$ , i.e. just a factorization algebra.

Poisson degree	???
1	$E_0$
0	$E_1$
-1	$E_2$
-2	$E_3$

**Definition I.2.** The  $BV_0$ -operad is the operad over  $\mathbb{R}[[\hbar]]$ , generated by a commutative product  $*$  and a Poisson bracket of degree 1, s.th.  $d* = \hbar\{\cdot, \cdot\}$ . ■

Invert  $\hbar$ :  $BV_0 \simeq E_0$ . If  $\hbar = 0$ , then  $BV_0/\hbar$  is the operad of commutative algebras with bracket of degree 1.

Let  $X$  be a manifold,  $f \in \mathcal{O}(X)$ . Then  $h(r, t(f))$ , the derived critical scheme has a  $\{\cdot, \cdot\}$  of degree 1 and "wants to become"  $E_0$ .  $\mathcal{O}(h(r, t(f)))$  is a dga

$$\Lambda^2 TX \xrightarrow{\iota_v(df)} TX \xrightarrow{\iota_v(df)} \mathcal{O}(X)$$

The Schouten bracket gives  $\mathcal{O}(h(r, t(f)))$  a Poisson bracket of degree 1. Apply this to the Euler-Lagrange situation and thus look at a "derived space of solutions". This is  $\text{Crit}(S)$  and should acquire the Poisson bracket.

**Example I.3.** In free field theory, the derived critical locus of  $\int \varphi \Delta \varphi$  is the 2-term complex

$$C^\infty(M) \xrightarrow{\Delta} C^\infty(M)$$

If we had an action with a cubic term, then differentiation acquires a non-linear term. ■

If  $B \subseteq M$ , the  $\mathcal{O}(hEL(B))$  looks like

$$\prod \text{Hom}(\text{Sym}^n(C^\infty(B)) \otimes \Lambda^k C^\infty(B), \mathbb{R})$$

with a differential coming from the action  $S$ . This has a bracket of degree 1, so it "wants to become"  $E_0$ , thus a factorization algebra. If  $X$  has a measure, i.e. a trivialization of the top-degree sheaf of differential forms, then we get a  $BV$  operator  $\Delta : \Lambda^\bullet TX \rightarrow \Lambda^\bullet TX$  such that  $d + \hbar \Delta$  gives an algebra over  $BV_0$ .

**Theorem I.4.** *The commutative factorization algebra, associated to the Euler-Lagrange equation of a classical field theory can be quantized into a factorization algebra in many interesting situations, e.g.*

- scalar field theories
- Yang-Mills theories on  $\mathbb{R}^4$ .

**Remark I.5.** This is false over  $\overline{\mathbb{Q}}$ ! ■