# Exotic Einstein metrics on $S^6$ and $S^3 \times S^3$ , nearly Kähler 6-manifolds and $G_2$ holonomy cones

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### G<sub>2</sub> cones and nearly Kähler 6–manifolds

Riemannian cone over smooth compact Riemannian manifold *M*:  $C(M) = \mathbb{R}^{+} \times M \text{ endowed with the Riemannian metric } g_{c} = dr^{2} + r^{2}g$ Hol(*C*)  $\subset$  G<sub>2</sub>  $\iff$  parallel (and hence closed) 3-form  $\varphi$  and 4-form  $*\varphi$   $\varphi = r^{2}dr \wedge \omega + r^{3}\text{Re}\Omega, \qquad *\varphi = -r^{3}dr \wedge \text{Im}\Omega + \frac{1}{2}r^{4}\omega^{2}$   $d\varphi = 0 = d * \varphi \iff \text{the SU(3)-structure } (\omega, \Omega) \text{ on } M \text{ satisfies}$  $\begin{cases} d\omega = 3 \text{Re}\Omega\\ d\text{Im}\Omega = -2\omega^{2} \end{cases}$ (NK)

A 6-manifold M endowed with an SU(3)-structure satisfying (NK) is called a (strict) **nearly Kähler** (nK) 6-manifold.

- every nK 6-manifold M is Einstein with Scal = 30  $\implies$  if M is complete, then it is compact with  $|\pi_1(M)| < \infty \implies$  wlog can assume  $\pi_1(M) = 0$ .
- nK 6-manifolds and real Killing spinors
- nK 2n-manifolds and Gray–Hervella classes of almost Hermitian manifolds

### The 4 examples known!

•  $S^6 \subset \operatorname{Im} \mathbb{O}$ : dates back to at least **1947** (e.g. C. Ehresmann, A. Kirchoff)

**1968**, in Gray–Wolf's classification of 3–symmetric spaces in 6d have

 $S^3 \times S^3 = SU(2)^3 / \triangle SU(2)$   $CP^3 = Sp(2) / U(1) \times Sp(1)$   $F_3 = SU(3) / T^2$ 

A 3-symmetric space has an automorphism  $\sigma$  with  $\sigma^3 = 1$ : define a homogeneous almost complex structure on ker ( $\sigma^2 + \sigma + Id$ ) by

$$J = \frac{1}{\sqrt{3}} \left( 2\sigma + \mathsf{Id} \right)$$

- Connection with G<sub>2</sub>-holonomy noted only in the 1980's, e.g. Bryant's 1987 first explicit example of a full G<sub>2</sub>-holonomy metric is C(F<sub>3</sub>)
- $G_2$ -cones give local models for isolated singularities of  $G_2$ -spaces
- Infinitely many Calabi–Yau, hyperkähler and Spin(7)–cones. Why not G<sub>2</sub>?
- 2005, Butruille: the four known examples are the only homogeneous nK 6-manifolds
- 2006, Bryant: local generality (via Cartan-Kähler theory) of 6d nK structures same as for 6d Calabi–Yaus (also Reves Carrion thesis 1993)

# Main Theorem and possible proof strategies

**Main Theorem** (Foscolo–Haskins, to appear **Annals of Mathematics**) There exists a complete inhomogeneous nearly Kähler structure on  $S^6$  and on  $S^3 \times S^3$ .

Two natural strategies to find nK 6-manifolds:

- Symmetries: cohomogeneity one nK 6-manifolds.
- Desingularisation of singular nK spaces.

Our proof uses elements from **both** viewpoints.

Simplest singular nK spaces: sine-cones (reduced holonomy  $SU(3) \subset G_2$ ) cross-section of a "split"  $G_2$  cone, i.e.  $\mathbb{R} \times C$  for C a Calabi–Yau cone  $(N^5, g_N)$  smooth Sasaki–Einstein  $\Leftrightarrow C(N)$  is a Calabi–Yau (CY) cone

The **sine-cone** over *N*:  $SC(N) = [0, \pi] \times N$  endowed with the Riemannian metric  $dr^2 + \sin^2 r g_N$  (aka **metric suspension** of *N*)

SC(N) is nK but has 2 isolated singularities each modelled on CY cone C(N)**Idea:** Try to desingularise SC(N) by replacing conical singularities with smooth asymptotically conical CY 3–folds.

### A simple nK sine-cone and desingularisations

A simple example comes from the so-called conifold:

- C(N) is the **conifold**  $\{z_1^2 + z_2^2 + z_3^2 + z_4^2 = 0\} \subset \mathbb{C}^4$
- $N = SU(2) \times SU(2) / \triangle U(1)$  which is diffeomorphic to  $S^2 \times S^3$

C(N) has 2 Calabi–Yau desingularisations (Candelas–de la Ossa, Stenzel)

- Y = the small resolution ≃ total space of O(-1) ⊕ O(-1) → P<sup>1</sup> vertex of cone replaced with a totally geodesic holomorphic P<sup>1</sup>
- Y' =the smoothing  $\simeq T^*S^3$

vertex of cone replaced with a totally geodesic special Lagrangian  $S^3$ The **conifold** itself and its asymptotically conical CY desingularisations are **cohomogeneity one**, i.e.  $\exists$  some Lie group *G* acting isometrically with generic orbit of codimension one

Two examples above have only 1 singular orbit:  $\mathbb{P}^1$  or  $S^3$ 

Sine-cone C(N), conifold and its desingularisations are cohomogeneity one. So obvious question is: **Can we desingularise this sine-cone as a cohomogeneity one space?** 

### Cohomogeneity one nK 6-manifolds

2010, **Podestà–Spiro**: potential complete cohomogeneity one nK 6–mfds M. Compact Lie group G acts with K,  $K_1$ ,  $K_2$  as its principal and singular isotropy groups. Principal orbit is G/K; 2 singular orbits  $G/K_i$ .

G	K	$K_1$	<i>K</i> <sub>2</sub>	Μ
SU(2)  imes SU(2)	riangle U(1)	riangleSU(2)	riangleSU(2)	$S^3  imes S^3$
$SU(2) \times SU(2)$	riangle U(1)	riangleSU(2)	U(1)  imes SU(2)	S <sup>6</sup>
$SU(2) \times SU(2)$	riangle U(1)	U(1)  imes SU(2)	SU(2)  imes U(1)	CP <sup>3</sup>
$SU(2) \times SU(2)$	riangle U(1)	U(1)  imes SU(2)	${\sf U}(1) imes{\sf SU}(2)$	$S^2  imes S^4$
SU(3)	SU(2)	SU(3)	SU(3)	$S^6$
$\Rightarrow$ $\textit{N}_{1,1}=\textit{SU}(2)\times\textit{SU}(2)/\Delta\textit{U}(1)$ is only possible interesting principal orbit!				

# Rough outline of proof

- 1. Understand the local theory for cohomogeneity one nK 6–mfds in neighbourhood of principal orbit  $N_{1,1} = SU(2) \times SU(2)/\Delta U(1)$ .
  - □ Our approach: study the geometry induced on (invariant) hypersurfaces and how it varies. Decomposes into a "static" and "dynamic" part.
  - □ Static = understand exactly what geometric structures can appear on an (invariant) hypersurface.

Answer = (invariant) *nearly hypo* SU(2) structures (Fernandez et al);

Space of invariant nearly hypo structures can be identified with a connected open subset of  $SO_0(1,2) \times S^1$ .  $S^1$  factor corresponds to obvious continuous symmetries of the equations.

So up to symmetry there exists a 3-dimensional family of invariant nearly hypo structures.

 Dynamic = (cohom 1) nK metrics correspond to differential equations for evolving a 1-parameter family of (invariant) nearly hypo structures.
Answer in cohom 1 case = explicit 1st order ODEs for a curve in the space of invariant nearly hypo structures.

 $\hfill\square$  Upshot:  $\exists$  2–parameter family of cohomogeneity 1 local nK metrics.

# Rough outline of proof II

• Don't know how to find explicit form for general solution to the ODEs. Special explicit solutions do exist, have geometric significance and play important role in our proof.

• Generic solution in 2-parameter family does NOT extend to a complete metric.

**Fundamental difficulty:** recognise which local solutions extend to complete metrics.

Proceed in two steps; separate the two singular orbits that appear and study separately.

**1.** Understand the possible singular orbits (uses Lie group theory) and which solutions extend over a given singular orbit (need to solve singular IVP).

2. Understand how to "match" a pair of solutions from the previous step. Step 1 fits into a general framework for cohomogeneity 1 Einstein metrics (Eschenburg–Wang 2000); extra care needed because of isotropy repn. Step 2 is the most subtle part of argument. Closest to previous work of Böhm on Einstein metrics on spheres (Inventiones 1998).

### Local solutions extending over a singular orbit

Neighbourhood of singular orbit is a *G*-equivariant disc bundle over singular orbit. Use representation theory to express conditions that a *G*-invariant section extend smoothly over the zero section. Get a *singular initial value problem* for 1st order nonlinear ODE system. Smoothness gives constraints on the initial values permitted.

Podestà-Spiro: up to symmetries possible singular orbits are

 ${\rm SU}(2) imes {\rm SU}(2)/{\rm U}(1) imes {\rm SU}(2) \simeq S^2$   ${\rm SU}(2) imes {\rm SU}(2)/ riangle {\rm SU}(2) \simeq S^3$ 

Proposition (Nearly Kähler deformations of small resolution & smoothing)

- There exist two 1-parameter families {Ψ<sub>a</sub>}<sub>a>0</sub> and {Ψ<sub>b</sub>}<sub>b>0</sub> of solutions to the fundamental ODE system which extend smoothly over a singular orbit S<sup>2</sup> and S<sup>3</sup>, respectively. a and b measure size of singular orbits.
- As a, b → 0, appropriately rescaled, the local nK structures Ψ<sub>a</sub> and Ψ<sub>b</sub> converge to the CY structures on the small resolution and the smoothing. Think of the two 1-parameter families as *local* nearly Kähler deformations of CY metrics on small resolution and smoothing.

Now the parameter a or b is NOT just a global rescaling (as in CY case).

# Matching pairs of solns: maximal volume orbits

*M* complete cohom 1 nK  $\implies$  orbital volume V(t) has a unique maximum. But generic member of our 1-parameter families of solutions is **not** complete.

Key properties of space of invariant maximal volume orbits  $\mathcal{V}$ :

- $\mathcal{V} \simeq \mathbb{R}^2 \times S^1 \subset \mathbb{R}^3 \times S^1$
- $V \ge 1$  on V and V = 1 precisely for the Sasaki–Einstein structure on N<sub>1,1</sub>
- $\mathcal{V} \cap \{V \leq C\}$  is compact

**Key Proposition** Every member of the families  $\{\Psi_a\}_{a>0}$  and  $\{\Psi_b\}_{b>0}$  has a unique maximal volume orbit.

Idea of proof: a continuity argument in the parameter a or b.

Nonempty; open; closed.

Nonempty: 3 of 4 known homogeneous examples appear in these families; these clearly have max vol orbits.

Openness: easy using nondegeneracy conditions that are satisfied.

Closedness is main point: uses compactness of  $\mathcal{V} \cap \{V \leq C\}$  plus standard ODE theory and basic comparison theory.

**Strategy for finding complete nK metrics**: Match pairs of solutions in the two families across their maximal volume orbits using discrete symmetries.

- $\alpha, \beta$  continuous curves in  $\mathbb{R}^2 \simeq \mathcal{V}/S^1$  parametrising the maximal volume orbits of  $\{\Psi_a\}_{a>0}$  and  $\{\Psi_b\}_{b>0}$
- Discrete symmetries = reflections along the axes
- Matching means:

(i) curves  $\alpha,\,\beta$  must intersect (up to a discrete symmetry), or

- (ii) self-intersect, or
- (iii) intersect either axis.
- Intersection points with axes correspond to solutions with a special "doubling symmetry", i.e. ∃ a reflection that exchanges the two singular orbits (therefore are of same type and size).
- Intersection points of  $\alpha$  curve with axes give  $S^2 \times S^4$  or  $CP^3$ .
- Intersection points of  $\beta$  curve with either axis gives  $S^3 \times S^3$ .
- Intersection points of  $\alpha$  and  $\beta$  curves (up to action of reflection) gives  $S^6$ . To understand if there exist new complete cohomogeneity 1 nK metrics is equivalent to:

How many axis crossings/intersection/self-intersection points do the curves  $\alpha$  and  $\beta$  (and their images under the reflections) have?

### Geometry of the $\alpha$ and $\beta$ curves

Two obvious ways to get some information about the  $\alpha$  and  $\beta$  curves.

1. Standard nK metrics on  $CP^3$ ,  $S^6$  and  $S^3 \times S^3$  give points on these curves.  $CP^3$  and  $S^3 \times S^3$  give intersection points of  $\alpha$  and  $\beta$  curves with the axes;  $S^6$  gives an intersection point between  $\alpha$  and (reflection of)  $\beta$  curves.

#### What about the sine-cone?

 Study limits of α and β curves as the parameters a and b → 0. Desingularisation philosophy suggests: Ψ<sub>a</sub> and Ψ<sub>b</sub> should both converge to the sine-cone away from the two singular orbits. Max vol orbit in sine-cone is the origin in the plane ("rotated" SE structure). So expect that the α and β curves both limit to the origin. Need to prove:

**Proposition.** As  $a, b \rightarrow 0$   $\Psi_a$  and  $\Psi_b$  converge to the sine-cone over the standard Sasaki–Einstein structure on N<sub>1,1</sub>.

*Proof ingredients:* use convergence of "bubbles" to asymptotically conical CY structures; the *Böhm functional*  $\mathcal{B}$  for cohom 1 Einstein metrics; invariance of  $\mathcal{B}$  under rescaling and fact that it gives a power of *Vol* on a max vol orbit; rotated SE metric is the absolute min of *Vol* on all max vol orbits.

### Existence of the new metric on $S^3 \times S^3$

First look only for solutions obtained by "doubling" some  $\Psi_b$ .

**Idea:** exploit the convergence of  $\Psi_b$  as  $b \to 0$  to the sine-cone and the existence of the homogeneous nK metric on  $S^3 \times S^3$  (this has b = 1). Find a new nK metric "between" these two metrics, i.e. with 0 < b < 1. **Observation:** Can detect a doubled metric on  $S^3 \times S^3$  via condition that

$$v_0 = 0$$

on a max volume orbit where  $v_0$  is one component of nearly hypo structure. *ODE* system  $\Rightarrow$  Zeros of  $v_0$  are nondegenerate; count the number of them that occur before a max vol orbit: we call this C(b).

Key fact: C(b) is locally constant in b unless we hit a doubled metric. Idea of proof: b = 1 is standard  $S^3 \times S^3$  and we check C(b) = 1. To get a new cohomogeneity 1 metric on  $S^3 \times S^3$  it's enough to show  $C(b) \ge 2$  for b > 0 sufficiently small. This implies there is another doubled metric for some  $b \in (0, 1)$ .

**Need to prove:**  $C(b) \ge 2$  for b > 0 sufficiently small.

## $C(b) \ge 2$ for b > 0 sufficiently small.

Want to count zeroes of  $v_0$  before max vol orbit; by ODE system this is equivalent to counting *critical points of*  $u_0$  (another component of the nearly hypo structure)

 $u_0$  is a solution of the second order IVP

(\*)  $(\lambda u_0')' + 12\lambda u_0 = 0, \quad u_0(0) = 0, \quad u_0'(0) = 2b^2 > 0.$ 

Convergence of  $\Psi_b$  to sine-cone implies  $\lambda(t) \rightarrow \sin t$  as  $b \rightarrow 0$ . **Idea:** Compare  $u_0$  to a solution of the limiting equation

 $(\sin t\,\xi')'+12\sin t\,\xi=0.$ 

This is Legendre's equation with k = 3. There are explicit solutions:

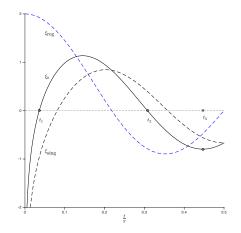
$$\xi_0(t) = C_1(5\cos^3 t - 3\cos t) + C_2\left(\frac{5}{2}\cos^2 t + \frac{1}{8}\cos t(4\cos^2 t - 6\sin^2 t)\log\frac{1 - \cos t}{1 + \cos t} - \frac{2}{3}\right)$$

1st solution is **regular** at endpoints 0 and  $\pi$  while 2nd is **singular**.

### Existence of the new metric on $S^3 \times S^3$

**Lemma:** There exists a solution  $\xi_0$  of this Legendre eqn with the following properties: there exists  $0 < t_1 < t_2 < t_3 < \frac{\pi}{2}$  such that  $\xi_0(t_1) = \xi_0(t_2) = 0$ ,  $\xi_0 \ge 0$  on  $[t_1, t_2]$  and  $\xi_0$  has a negative minimum at  $t_3$ .

Proof:



# Existence of the new metric on $S^3 \times S^3$

Recall that  $u_0$  solves

(\*)  $(\lambda u'_0)' + 12\lambda u_0 = 0$ ,  $u_0(0) = 0$ ,  $u'_0(0) = 2b^2 > 0$ .

and  $\lambda \rightarrow \sin t$  on  $(0, \pi)$  as  $b \rightarrow 0$ .

**Theorem:** There exists  $\epsilon > 0$  such that for all  $b < \epsilon$ ,  $u_0$  the solution of (\*) has a strict negative minimum before the maximal volume orbit.

*Proof sketch:* Apply a (generalised) Sturm-Picone comparison argument to prove the same conclusion about the minimum holds for solution of (\*), using uniform convergence of  $\lambda(t)$  to sin t on compact subsets of  $(0, \pi)$ . Finally: initial conditions for  $u_0$  also force a *maximum* before the minimum.

### **Existence of the new metric on** $S^6$

Need to force planar curves  $\alpha$  and  $\beta$  to intersect in another point (2 intersection points already exist: standard nK  $S^6$  and sine-cone).

**Idea:** use the new and old solutions on  $S^3 \times S^3$  to find a closed bounded region D in the plane encircling the origin. The  $\alpha$  curve starts at the origin (as  $a \rightarrow 0$ ); we want to show that eventually the  $\alpha$  curve leaves D passing through its boundary; this point gives the new intersection point of the  $\alpha$  and  $\beta$  curves.

**Proposition.** The curve  $\alpha$  exits any compact subset of  $\mathbb{R}^2$  as  $a \to \infty$ .

*Idea of proof:* based on explicit Taylor series for solutions  $\Psi_a$  and their dependence on the parameter *a* we consider a very particular (but non geometric) rescaling of the solution to the ODE system. Show that the rescaled solutions are well behaved as  $a \to \infty$  and converge smoothly to some limiting object. Scaling used shows  $V_{max} \sim ca^4$ .

*Heuristically:* making the size of the singular orbit 2-sphere large  $(a \rightarrow \infty)$  forces the size of the maximal volume orbit to be large.

### The intersection points of $\alpha$ and $\beta$

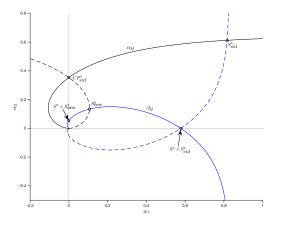


Figure:  $\alpha$  and  $\beta$  curves and the locations of the 5 complete cohomogeneity one nK structures computed numerically

### Nearly hypo structures

$$\begin{split} \mathsf{N}_{1,1} &= \mathsf{SU}(2) \times \mathsf{SU}(2) / \triangle \mathsf{U}(1) \\ \mathsf{On} \ M^* &= (a,b) \times \mathsf{N}_{1,1} \text{ we write} \end{split}$$

 $\omega = \eta \wedge dt + \omega_1$   $\Omega = (\omega_2 + i\omega_3) \wedge (\eta + idt),$ 

where  $(\eta, \omega_1, \omega_2, \omega_3)$  defines an invariant SU(2)-structure on N<sub>1,1</sub>.

The SU(2)-structure induced on a hypersurface in a nK 6-manifold is called a **nearly hypo** structure. It is defined by the following equations:

$$d\omega_1 = 3\eta \wedge \omega_2$$
  $d(\eta \wedge \omega_3) = -2\omega_1 \wedge \omega_1$ 

The evolution equations to obtain a nK manifold by flowing a nearly hypo structure are:

 $\partial_t \omega_1 = -d\eta - 3\omega_3$   $\partial_t (\eta \wedge \omega_2) = -d\omega_3$   $\partial_t (\eta \wedge \omega_3) = d\omega_2 + 4\eta \wedge \omega_1$ 

**Lemma.** The space of invariant nearly hypo structures on N<sub>1,1</sub> is a smooth manifold diffeomorphic to  $\mathbb{R}^3 \times S^1$ . (The S<sup>1</sup>-factor is generated by the action of the Reeb vector field.)

### The fundamental ODE system

We parametrise invariant nearly hypo structures modulo the Reeb action by tuples  $(\lambda, u, v) \in \mathbb{R}^+ \times \mathbb{R}^{1,2} \times \mathbb{R}^{1,2}$  subject to the constraints

$$\lambda^2 |u|^2 = |v|^2 > 0$$
  $\langle u, v \rangle = 0$   $v_1 = |u|^2$   $u_2 = -\lambda |u|$ 

The basic equations then are:

$$\begin{split} \lambda \dot{u}_0 + 3v_0 &= 0, & \dot{v}_0 - 4\lambda u_0 &= 0, \\ \lambda \dot{u}_1 + 3v_1 - 2\lambda^2 &= 0, & \dot{v}_1 - 4\lambda u_1 &= 0, \\ \lambda \dot{u}_2 + 3v_2 &= 0, & \lambda \dot{v}_2 - 4\lambda^2 u_2 + 3u_2 &= 0, \end{split}$$

$$\lambda^{2}|u|^{2}\dot{\lambda}+2\lambda^{4}u_{1}+3u_{2}v_{2}=0.$$

**Proposition.** Up to symmetries, there exists a 2-parameter family of local cohomogeneity one nK stuctures on  $(a, b) \times N_{1,1}$ .

The homogeneous nK structure on  $S^3 \times S^3$  is

$$\lambda = 1, \quad u_0 = u_1 = \frac{1}{\sqrt{3}} \sin(2\sqrt{3}t), \quad u_2 = -\frac{2}{\sqrt{3}} \sin(\sqrt{3}t),$$

$$\begin{aligned} v_0 &= -\frac{2}{3}\cos\left(2\sqrt{3}t\right), \quad v_1 = \frac{2}{3}\left(1 - \cos\left(2\sqrt{3}t\right)\right), \quad v_2 = \frac{2}{3}\cos\left(\sqrt{3}t\right), \end{aligned}$$
 for  $t \in [0, \frac{\pi}{\sqrt{3}}].$ 

The sine-cone is:

$$\lambda = \sin t$$
,  $u_0 = 0$ ,  $u_1 = \sin^2 t \cos t$ ,  $u_2 = -\sin^3 t$ ,

$$v_0 = 0$$
,  $v_1 = \mu^2 = \sin^4 t$ ,  $v_2 = \sin^3 t \cos t$ ,

for  $t \in [0, \pi]$ .

The first few terms of the Taylor series of  $\Psi_a$  at t = 0 are:

$$\lambda(t) = \frac{3}{2}t - \frac{2a^2 + 3}{12a^2}t^3 + \frac{116a^4 - 381a^2 + 261}{1440a^4}t^5 + \cdots$$
$$u_0(t) = a^2 - 3a^2t^2 + \frac{52a^2 - 3}{24}t^4 + \cdots$$

$$u_1(t) = a^2 - \frac{3}{2}(2a^2 - 1)t^2 + \frac{52a^4 - 32a^2 - 3}{24a^2}t^4 + \cdots$$

$$u_2(t) = rac{-3\sqrt{3}}{2}at^2 + rac{\sqrt{3}(16a^2 - 3)}{12a}t^4 + \cdots$$

# Conclusion

**Conjecture** The Main Theorem yields all (inhomogeneous) cohomogeneity one nK structures on simply connected 6–manifolds.