

*Topological Quantum Field Theories: Knots
and Links in Three-dimensions and Black
Holes in $3 + 1$ Dimensions*

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On the other hand, topological (global) properties are independent of the metric.

For example, the size and shape of a knot in a three dimensional manifold do depend on the metric, its 'knottedness' does not.

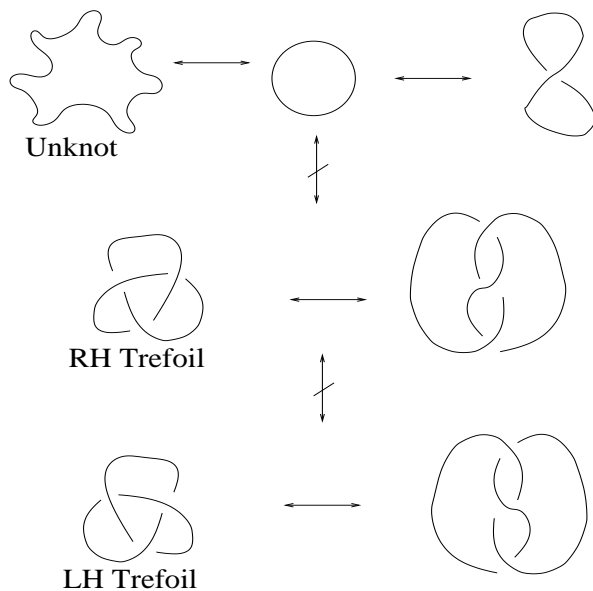
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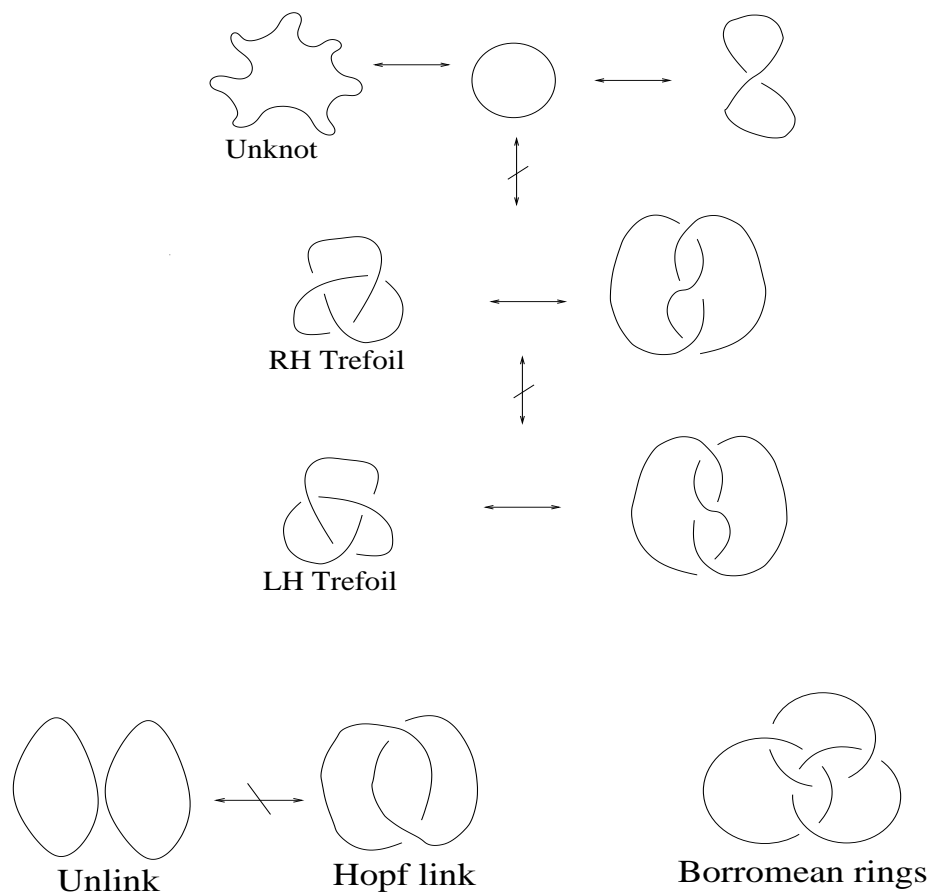
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Reviews: Birmingham *et al*, Topological field theory, Phys. Rep. 209 (1991),129.

Kaul, Govindarajan, Ramadevi, Schwarz type topological quantum field theories, in

Encyclopedia of Mathematical Physics, 2006, 494, hep-th/0504100.

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Early discussion of Hamiltonian quantization:

Bos and Nair: Phys. Letts. **B223**, (1989), 61; Int.J.Mod.Phys. **A5**, (1990), 959.

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Thus, for any compact gauge group \mathcal{G} , we have a new polynomial knot/link invariant.

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Artin Braids: Braids of unoriented single colour strands form a group.

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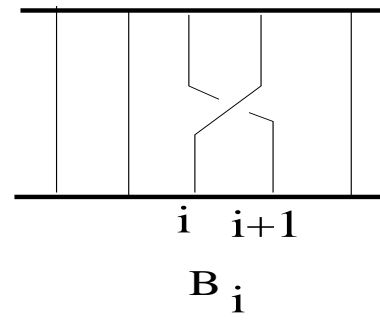
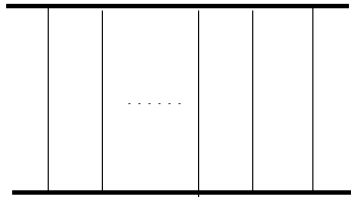
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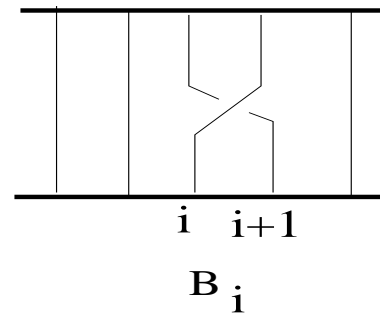
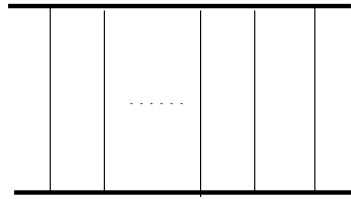
Construction of this complete and explicit solution exploits the close connection of knots/links with braids.

(RKK: Commun. Math.Phys. 162, 1994, 289.)

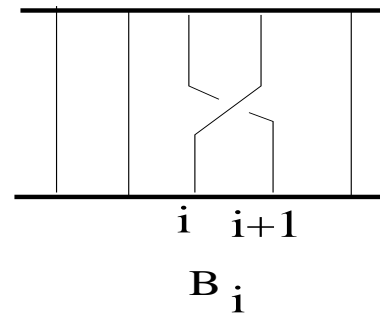
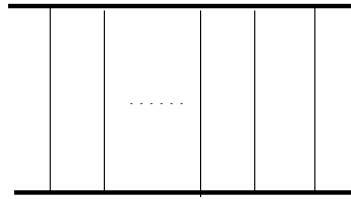
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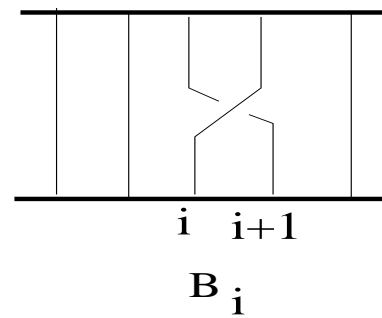
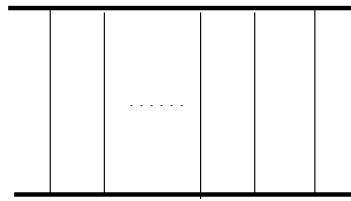


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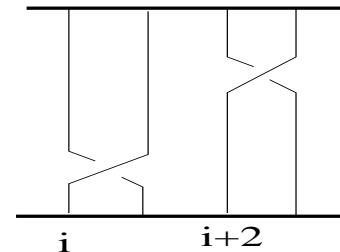
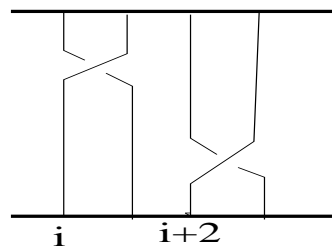
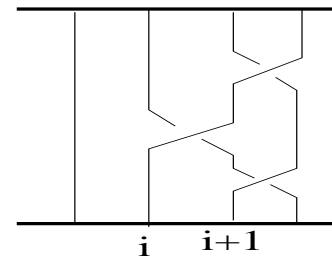
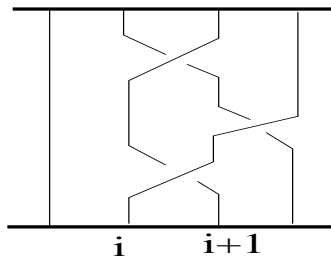
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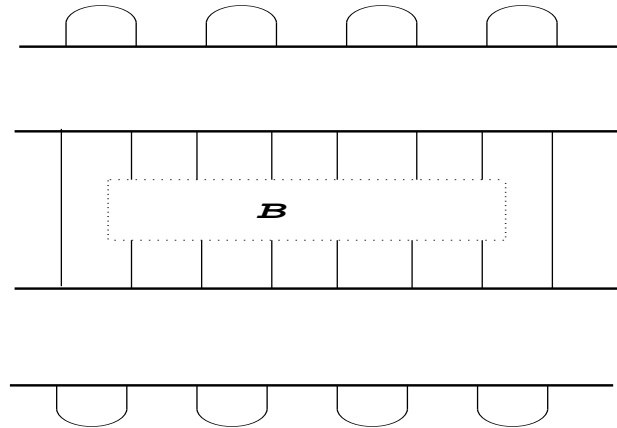
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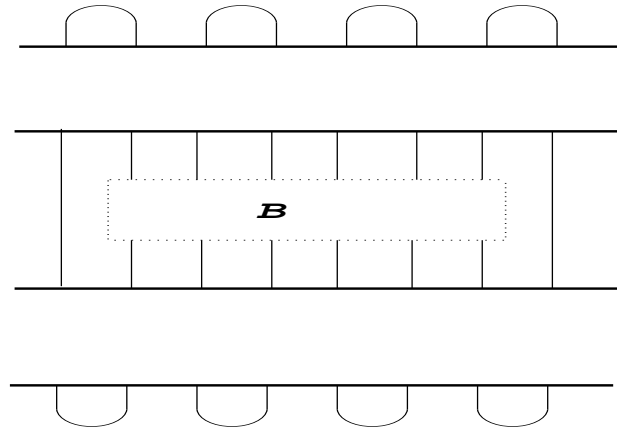
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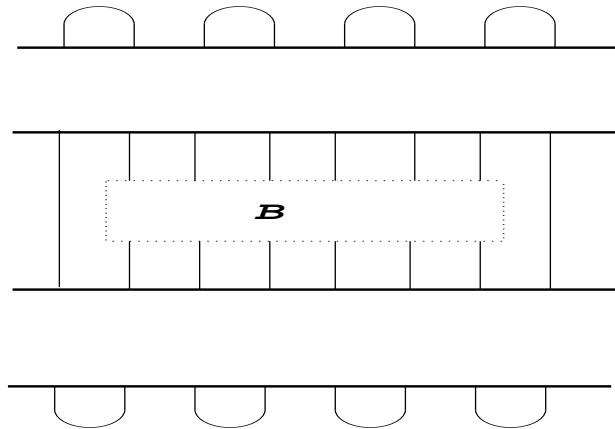
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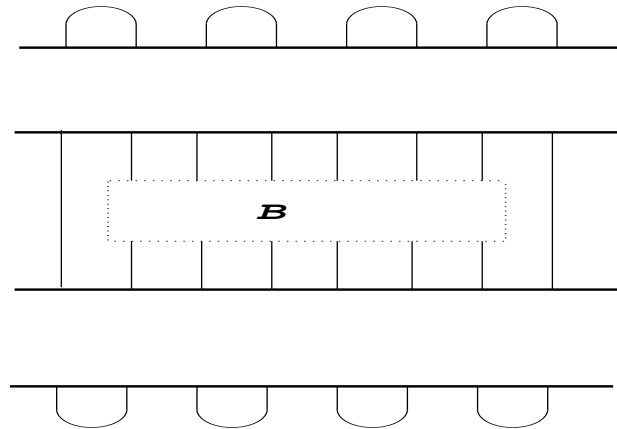


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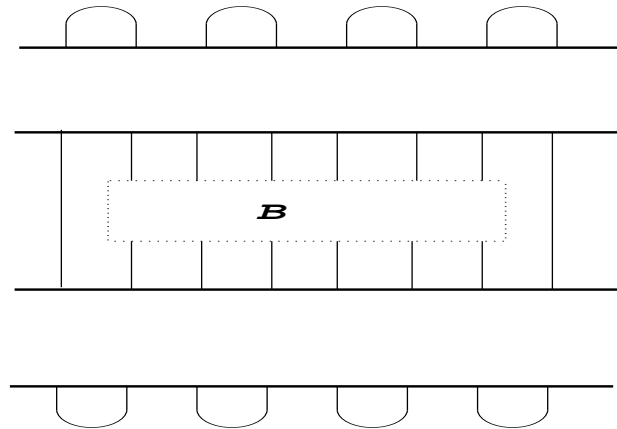
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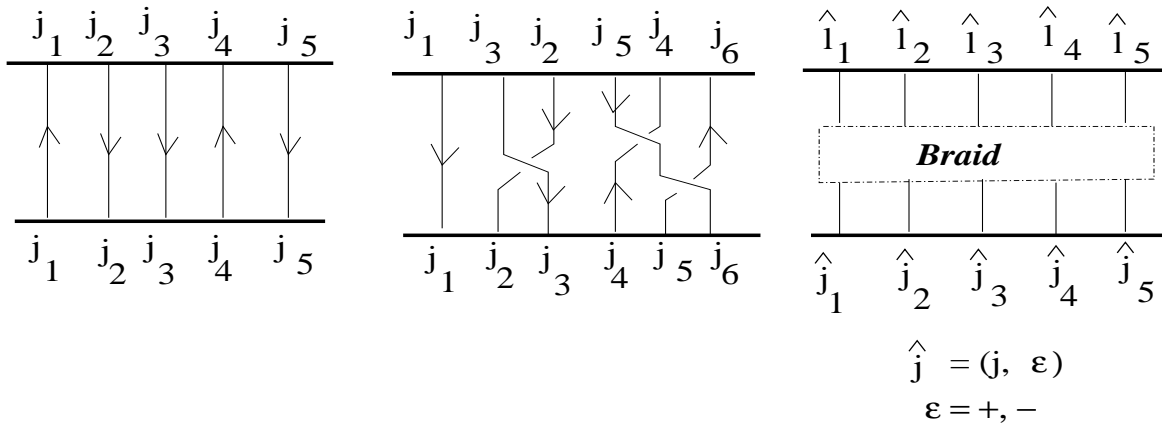
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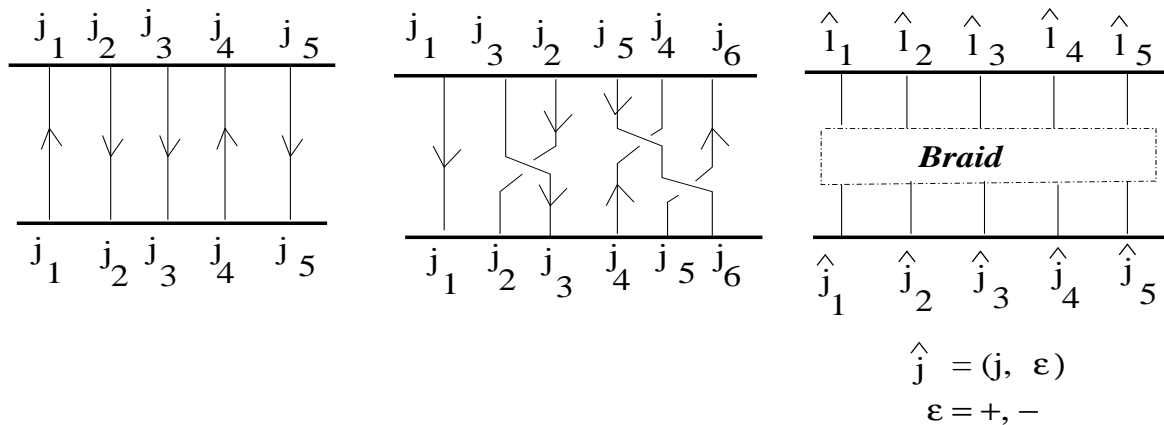


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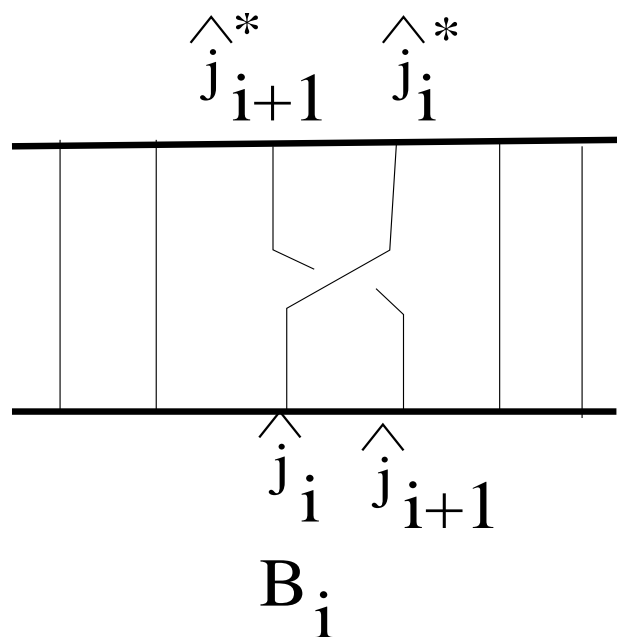
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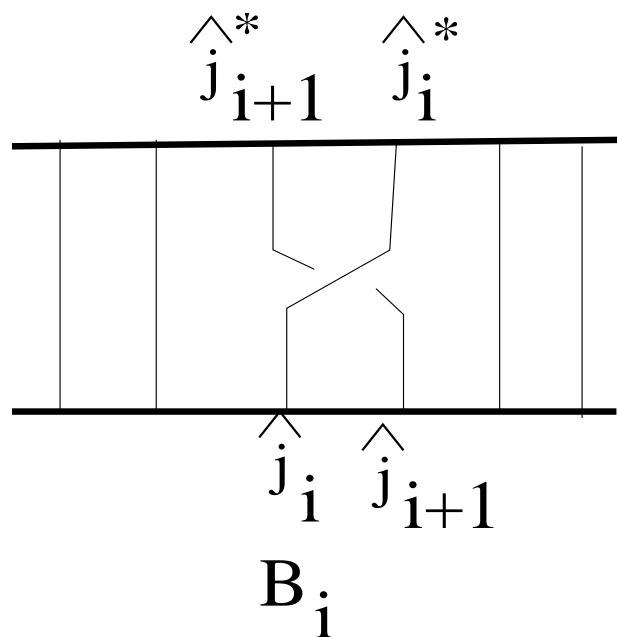
The assignments $(\hat{l}_1, \hat{l}_2, \hat{l}_3, \hat{l}_4, \hat{l}_5)$ are a permutation of $(\hat{j}_1^*, \hat{j}_2^*, \hat{j}_3^*, \hat{j}_4^*, \hat{j}_5^*)$, where $\hat{j}_i = (j_i, \pm)$ and $\hat{j}_i^* = (j_i, \mp)$.

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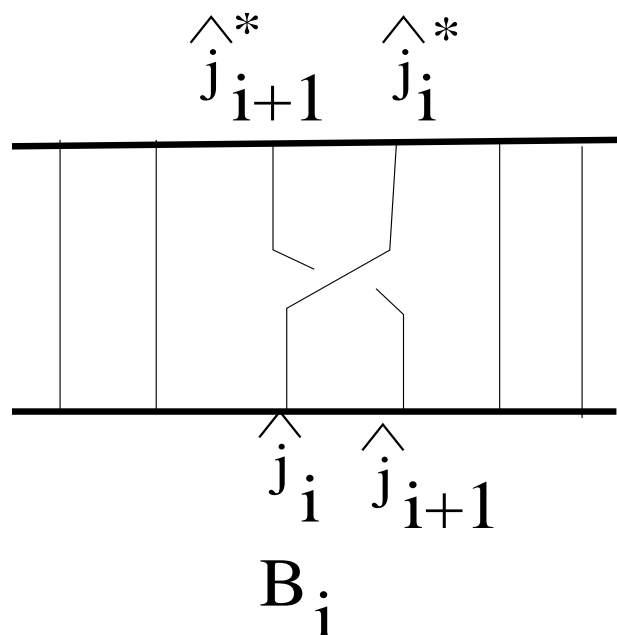


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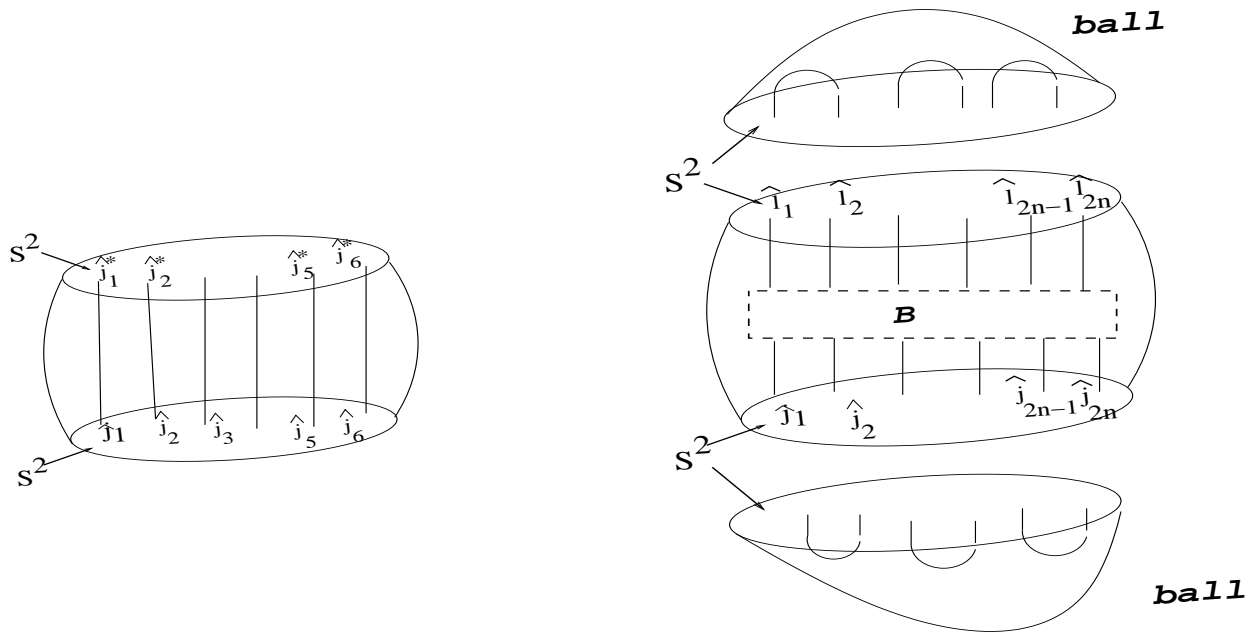
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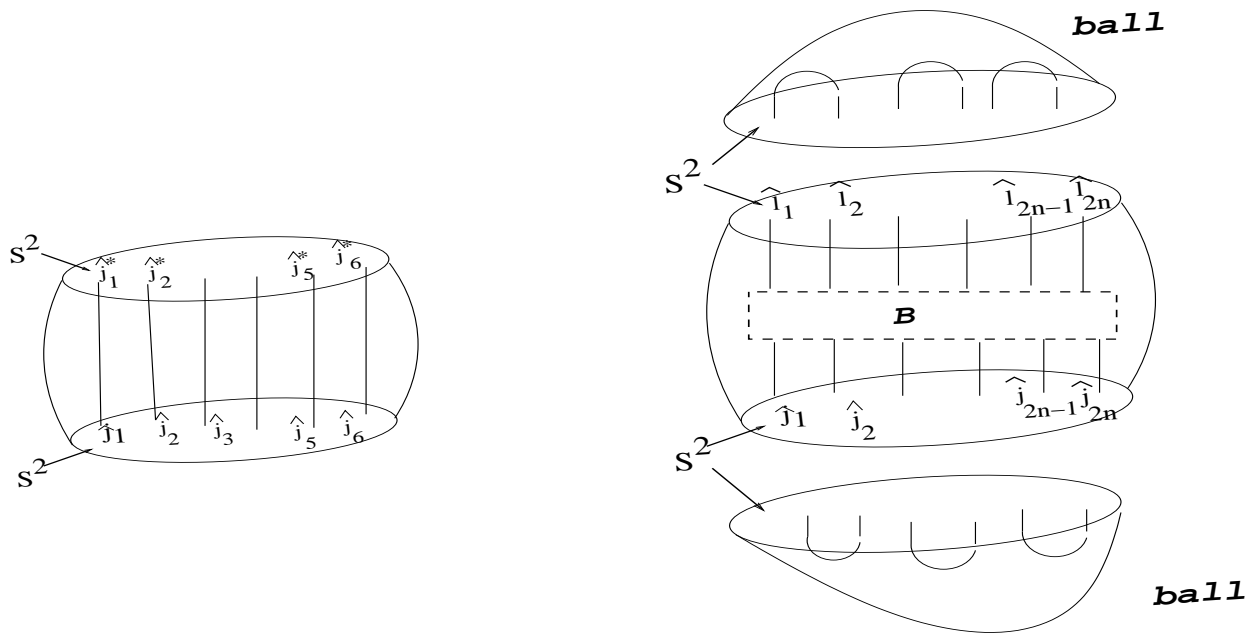
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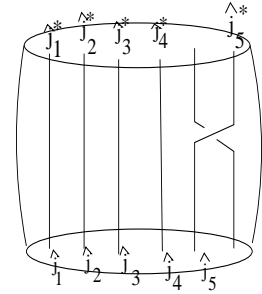
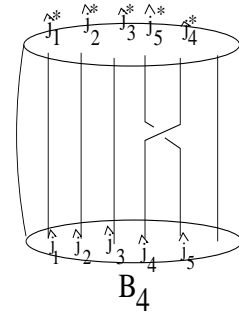
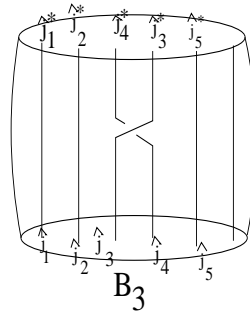
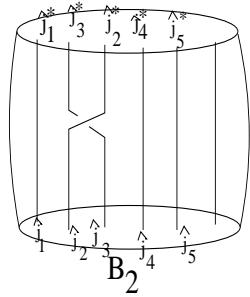
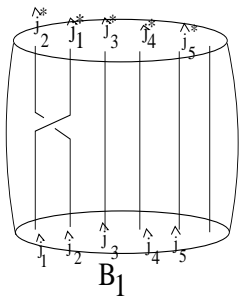
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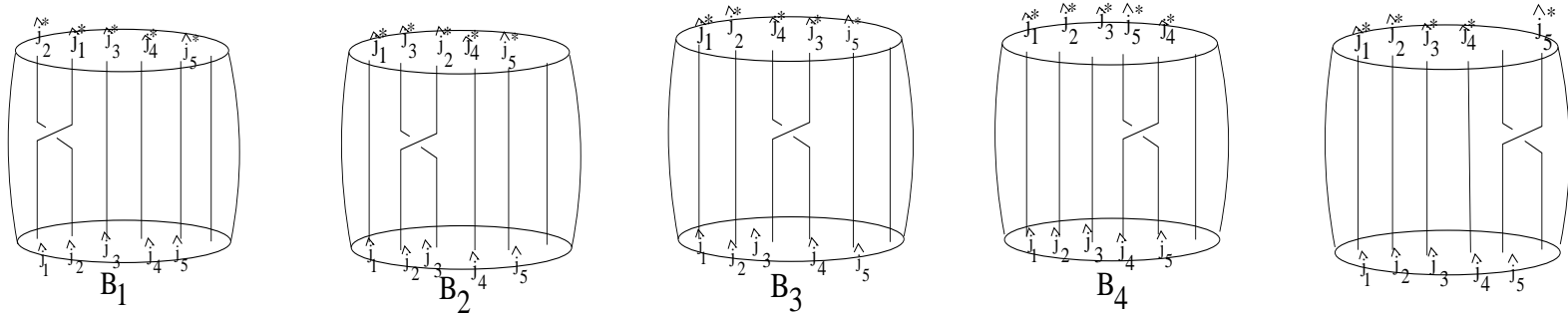
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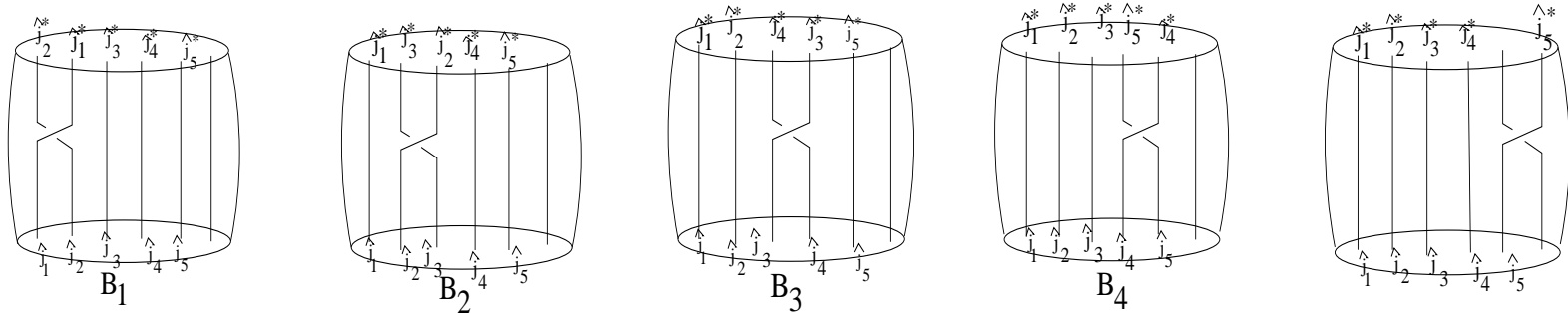


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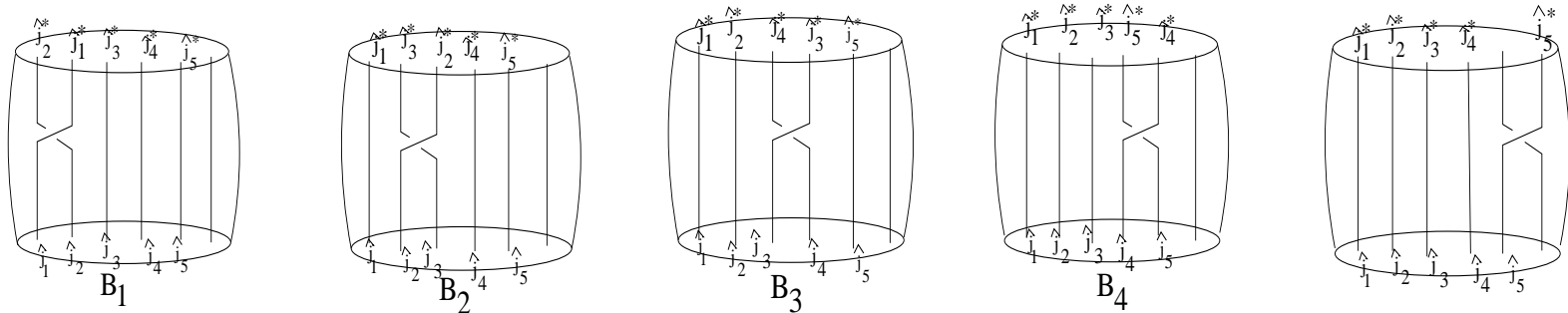
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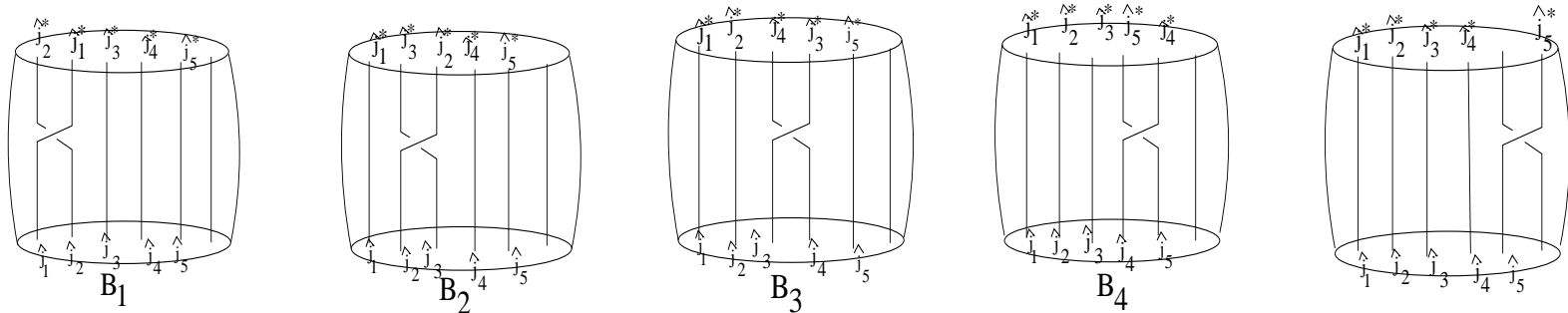


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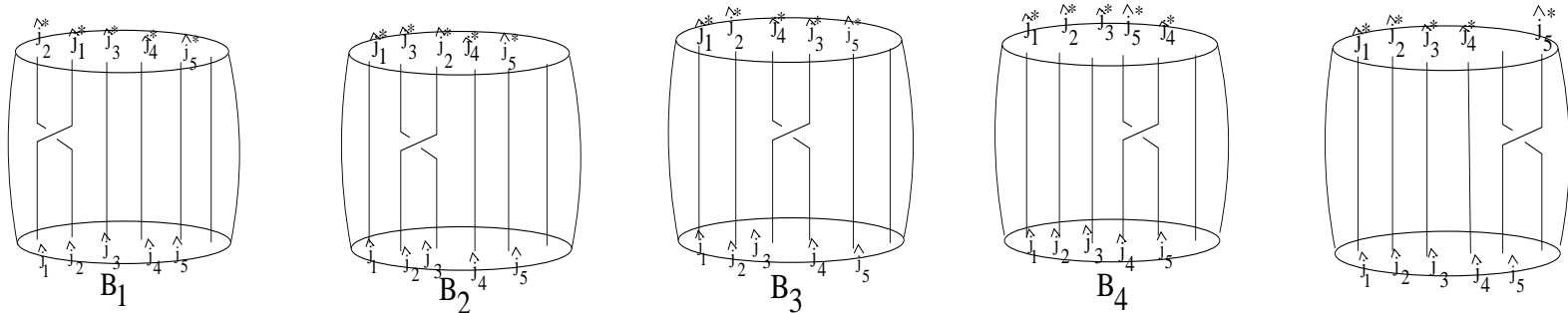
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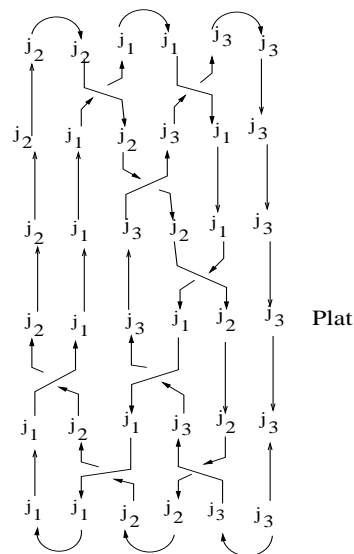
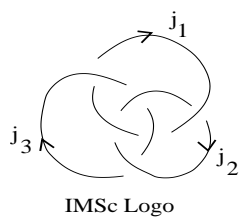
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$V_{j_1 j_2 j_3}$ [IMSc logo]

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$$a_{j\ell} \begin{bmatrix} j_1 & j_2 \\ j_3 & j_4 \end{bmatrix} = (-)^{(j_1+j_2+j_3+j_4)} \sqrt{[2j+1][2\ell+1]} \begin{pmatrix} j_1 & j_2 & j \\ j_3 & j_4 & \ell \end{pmatrix}_q, \quad [x] = \frac{q^{x/2} - q^{-x/2}}{q^{1/2} - q^{-1/2}}$$

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RKK: *Entropy of Quantum Black Holes*, SIGMA 8, (2012), 005.

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$$S_{CS} = \frac{k}{2\pi} \int_{\Delta} \epsilon^{abc} \left(A_a^i \partial_b A_c^i + \frac{1}{3} \epsilon^{ijk} A_a^i A_b^j A_c^k \right) + \int_{\Delta} J^{ai} A_a^i$$

where $J^{ai} = (J^i, 0, 0)$.

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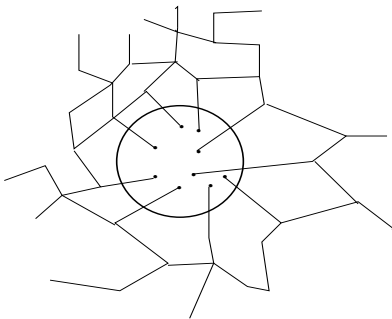
Black hole micro-states

Quantizing this CS theory then provides a framework to count the horizon micro-states.

Quantum theory is described by quantizing the $SU(2)$ CS theory in the bdy Δ with bulk solder form

$J^i = \Sigma_{\theta\phi}^i \equiv \frac{1}{2}\epsilon^{ijk} e_{\theta}^j e_{\phi}^k$ as the external source.

LQG is described by spin networks made from functionals of Wilson line integral as configuration operators. The corresponding momentum operators are the flux operators constructed from the solder forms: $\int_{S^2} d^2\sigma \Sigma_{\theta\phi}^i$.



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Now the micro-states can be counted by simply counting the number of states satisfying this constraint, either in bulk theory or in the boundary CS theory.

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$$\mathcal{N}_{\mathcal{P}}(j_1, j_2, \dots, j_p) = \frac{2}{k+2} \sum_{r=0}^{k/2} \frac{\prod_{l=1}^p \sin\left(\frac{(2j_l+1)(2r+1)\pi}{k+2}\right)}{\left[\sin\left(\frac{(2r+1)\pi}{k+2}\right)\right]^{p-2}}.$$

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$$A_{\text{H}} = 8\pi\gamma \sum_{l=1,2,\dots,p} \sqrt{j_l(j_l + 1)}; \quad \ell_P = 1.$$

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All these facts suggest a universality of this entropy formula.

Conclusion

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In the process a new class of matrix representations of coloured oriented braids are obtained, which have wide applications.

$SU(2)$ CS theory also provides a description of the black hole micro-states, yielding an entropy formula with a possibly universal coefficient, $-3/2$, for the logarithmic area correction to the Bekenstein-Hawking law.