Limitations of Geometric Engineering: Implications for Model Building

David R. Morrison

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Limitations of Geometric Engineering: [Implications for](#page-53-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \blacktriangleright In the geometric engineering approach, properties of the vacua which arise from local features of the geometry are studied purely locally.

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

 209

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \blacktriangleright Typically, one studies spacetimes of the form $X \times M$. where X is a noncompact geometric space containing the local feature in question, and M is Minkowski space (or sometimes de Sitter or anti de Sitter space).
- \triangleright My theme today: globalizing the local constructions can be delicate and dangerous.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

- \triangleright As a first example of the difficulties in globalizing, consider the ADE singularities.
- ► These can be studied locally as $X_0 = \mathbb{C}^2/\Gamma$, with their resolutions given by ALE spaces X . Much can be learned about the rôle which such singularities play in string vacua by studying $X_0 \times M$ or $X \times M$. The presence of the singularity leads to a nonabelian gauge group in IIA string theory, M-theory and F-theory.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

A O A G A 4 O A C A G A G A H 4 O A C

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- \blacktriangleright However, globalizing these singularities is a different matter. If one asks for a Calabi–Yau 2-fold (i.e., a K3 surface) with such a singularity, then the rank of the gauge group is bounded by 19.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

- \triangleright As a first example of the difficulties in globalizing, consider the ADE singularities.
- ► These can be studied locally as $X_0 = \mathbb{C}^2/\Gamma$, with their resolutions given by ALE spaces X . Much can be learned about the rôle which such singularities play in string vacua by studying $X_0 \times M$ or $X \times M$. The presence of the singularity leads to a nonabelian gauge group in IIA string theory, M-theory and F-theory.
- \blacktriangleright However, globalizing these singularities is a different matter. If one asks for a Calabi–Yau 2-fold (i.e., a K3 surface) with such a singularity, then the rank of the gauge group is bounded by 19.
- In this case, it is possible to explicitly describe which gauge groups can occur (cf. DRM, Invent. Math. 75 (1984) 105–121).

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0) [Conclusions](#page-49-0)

KORKAPRASHASH SA MAG

 \triangleright Suppose one wanted to do some kind of statistical analysis of gauge groups in this context. Based on the local analysis, one would likely conclude that groups of type $SU(n)$ or $SO(2n)$ occur with probability 1, and the groups of type E_6 , E_7 , and E_8 occur with probability 0.

Geometric Engineering: [Implications for](#page-0-0) Model Building

Limitations of

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright The global story is much different: because the ranks are bounded, the E_n groups play a significant rôle in the overall statistics, and occur with probability greater than 0.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0) [Conclusions](#page-49-0)

A O A G A 4 O A C A G A G A H 4 O A C

7-branes in F-theory

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright The essential data for an F-theory vacuum is the (complexified) IIB coupling, which is determined from the j-invariant of the family of elliptic curves

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y^2 = x^3 + px + q,
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where p and q are functions on the space X being used to (partially) compactify the IIB string.

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

 209

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 \triangleright The 7-brane wraps the discriminant locus $\{\Delta = 0\}$ of the above equation, where

$$
\Delta = 4p^3 + 27q^2.
$$

When X is local, we can assume that p and q have no common zeros.**K ロ ▶ K 個 ▶ K 할 ▶ K 할 ▶ 그럴 → 이익단**

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

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 \triangleright But in the global case, if X has complex dimension at least 2, p and q are forced to have common zeros (except in very special cases, like product manifolds).

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

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- At a common zero of p and q, $\{\Delta = 0\}$ is itself singular (it has a "cusp" singularity) and the standard analysis of a 7-brane as a submanifold is not correct.
- In fact, there are a number of global issues with 7-branes in F-theory; in six dimensions, many of them can be related to anomaly cancellation in the effective six-dimensional theory (cf. Grassi–DRM, math.AG/0005196).

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

 \triangleright Note that F-theory itself provides an alternative to pure D-brane constructions in this context.

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

 \triangleright Note that F-theory itself provides an alternative to pure D-brane constructions in this context. The global data needed for F-theory vacua is specified by the family of elliptic curves (together with a bit more data pertaining to bundles on the branes), and methods of algebraic geometry can be used to probe the structure of these models in detail (as discussed in other talks at this conference).

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

A O A G A 4 O A C A G A G A H 4 O A C

 \triangleright String vacua with D-branes (and orientifold planes) are ubiquitous in modern string phenomenology. But the global issues in the presence of D-branes are particularly subtle.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- If the D-brane charges do not cancel, leaving a net Ramond–Ramond charge in the vacuum, there is a tadpole anomaly.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

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- \triangleright This tadpole is suppressed in local models, since any excess Ramond–Ramond charge can escape to infinity. But it cannot be ignored in global models.

A O A G A 4 O A C A G A G A H 4 O A C

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

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- If the D-brane charges do not cancel, leaving a net Ramond–Ramond charge in the vacuum, there is a tadpole anomaly.
- \triangleright This tadpole is suppressed in local models, since any excess Ramond–Ramond charge can escape to infinity. But it cannot be ignored in global models.
- \triangleright Still, it has seemed reasonable to study different sources of D-brane charge in a local way, and assemble them later into a global model. This approach has recently been called into question.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

 \triangleright Typically, much of our information about local models comes not directly from analyzing string theory itself, but by twisting to get a topological theory and then studying that topological theory locally.

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

- \triangleright Typically, much of our information about local models comes not directly from analyzing string theory itself, but by twisting to get a topological theory and then studying that topological theory locally.
- Recently, in work of Walcher, arXiv:0712.2775 [hep-th], and Cook–Ooguri–Yang, arXiv:0804.1120 [hep-th], it was realized that beyond tree-level, the open topological string theory itself is sensitive to the presence of a net Ramond–Ramond charge, when studied in the context of global models.

A O A G A 4 O A C A G A G A H 4 O A C

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

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- Recently, in work of Walcher, $arXiv:0712.2775$ [hep-th], and Cook–Ooguri–Yang, arXiv:0804.1120 [hep-th], it was realized that beyond tree-level, the open topological string theory itself is sensitive to the presence of a net Ramond–Ramond charge, when studied in the context of global models.
- \triangleright That is, if one attempts to calculate open topological string amplitudes in a background with net Ramond–Ramond charge (using techniques pioneered by Walcher a few years ago), the answers reveal a new kind of anomaly in the topological string.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

 \triangleright This suggests that the purely local contributions to topological string amplitudes which have been computed in many situations will require unexpected corrections from other sectors of the background. It is not clear at present which conclusions about local physics escaped unchanged.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

- \triangleright This suggests that the purely local contributions to topological string amplitudes which have been computed in many situations will require unexpected corrections from other sectors of the background. It is not clear at present which conclusions about local physics escaped unchanged.
- \blacktriangleright In particular, model building which has relied on a decoupling of the D-brane and orientifold plane sectors of the theory may need to be reexamined.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0) [Conclusions](#page-49-0)

A O A G A 4 O A C A G A G A H 4 O A C

Conifold transitions

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright The local analysis starts with a 3-cycle in a local Calabi–Yau threefold X which is the vanishing cycle for an ordinary double point singularity (when the complex structure parameters are specialized), and the IIA string on $X \times M$ with a D6-brane wrapping the 3-cycle N times. Passing to the double point, and then blowing it up to obtain a 2-sphere, one finds a flux vacuum with N units of Ramond–Ramond flux on the 2-sphere.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

KORKAPRASHASH SA MAG

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- \blacktriangleright This transition has a lift to M-theory which was studied by Atiyah–Maldacena–Vafa, hep-th/0011256, Atiyah–Witten, hep-th/0107177, and others. The local M-theory space Y is a neighborhood of $S^3 \times S^3$, and the open string conifold transition is seen as a smooth deformation of M-theory vacu[a.](#page-30-0) of the service of the service

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

 \blacktriangleright In order to realize a conifold transition on a compact Calabi–Yau manifold, the collection of 3-cycles on which the transition is being made must have a homology relation (cf. Greene–DRM–Strominger, hep-th/9504145).

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- In the context of the open string model, if we use the coefficients of the homology relation as the N_i on the various 3-cycles, then the net D6-brane charge is zero! Thus, the global requirement from geometry (in order to get a Kähler manifold after the transition) takes care of the tadpole anomaly in string theory as well as the open string anomaly discussed above.

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- In the context of the open string model, if we use the coefficients of the homology relation as the N_i on the various 3-cycles, then the net D6-brane charge is zero! Thus, the global requirement from geometry (in order to get a Kähler manifold after the transition) takes care of the tadpole anomaly in string theory as well as the open string anomaly discussed above.
- \blacktriangleright How does this lift to M-theory? (cf. DRM, to appear).

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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 \triangleright Given a compact Calabi–Yau manifold X and a collection of 3-cycles Z , the lift to M-theory should be a G2-manifold Y with a circle action that has fixed points along the 3-cycles $Z \subset Y$.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright At present, unlike the local case, we do not have existence theorems for such G2-manifolds. However, the topological structure (and hence the classical moduli space) of our postulated G2-manifold can be analyzed. In general, one should consider both fixed points of a circle action, and allow the situation in which the Chern class of the circle bundle over $X - Z$ is non-trivial. (These are the flux vacua.)

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric **[Engineering](#page-1-0)**

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0) [Conclusions](#page-49-0)

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- \blacktriangleright The result is that the G2-moduli space has an extra parameter for each 3-cycle in Z and an extra parameter for each 2-cycle in the support of the Chern class of the circle bundle.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

 $\mathbf{E} = \mathbf{A} \mathbf{E} \mathbf{b} + \mathbf{A} \mathbf{E} \mathbf{b} + \mathbf{A} \mathbf{B} \mathbf{b} + \mathbf{A} \mathbf{b}$ 209

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- \blacktriangleright (These are "extra" when compared with the usual calculation of moduli for the case $X\times S^1.$) These extra parameters allow for the Atiyah et al. explanation of the local conifold transition to be present in the global moduli space.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright We have in mind the class of models which involve both D6-branes and O6-planes. (Global consistency for other types of D-brane charge will have to be addressed in a different way.) The lift of the D6-brane has already been discussed: it is a codimension 4 fixed point locus of a circle action, which has a metric that is asymptotically the Taub–NUT metric. The lift of the O6-plane is also known: the metric in this case is asymptotically the Atiyah–Hitchin metric which describes the two-monopole solution in Yang–Mills theory (cf. Seiberg–Witten, hep-th/9607163).

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \blacktriangleright The expectation from the past half-decade of string phenomenological studies is that there will be a large number of such vacua, and hence a large number of G2-manifolds which lift them to M-theory.
- \triangleright Do these all exist? At the moment, the best we can do is to analyze topological restrictions on such manifolds. As indicated above, I've already done this for cases without orientifold planes; work on the more general case is in progress.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright plugging those features into global models can be tricky.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)

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- \triangleright Even the local analysis for an open string model, done using topological string theory, can be misleading.
- \triangleright Global approaches, such as F-theory instead of orientifolds of IIB, or lifing to M-theory instead of orientifolds of IIA, hold much promise for resolving global issues.

Limitations of Geometric Engineering: [Implications for](#page-0-0) Model Building

David R. Morrison

Geometric [Engineering](#page-1-0)

[ADE singularities](#page-5-0)

[7-branes in](#page-11-0) F-theory

[Open strings and](#page-20-0) D-branes

Conifold [transitions](#page-29-0)

[Lifting to M-theory](#page-41-0)

[Conclusions](#page-49-0)